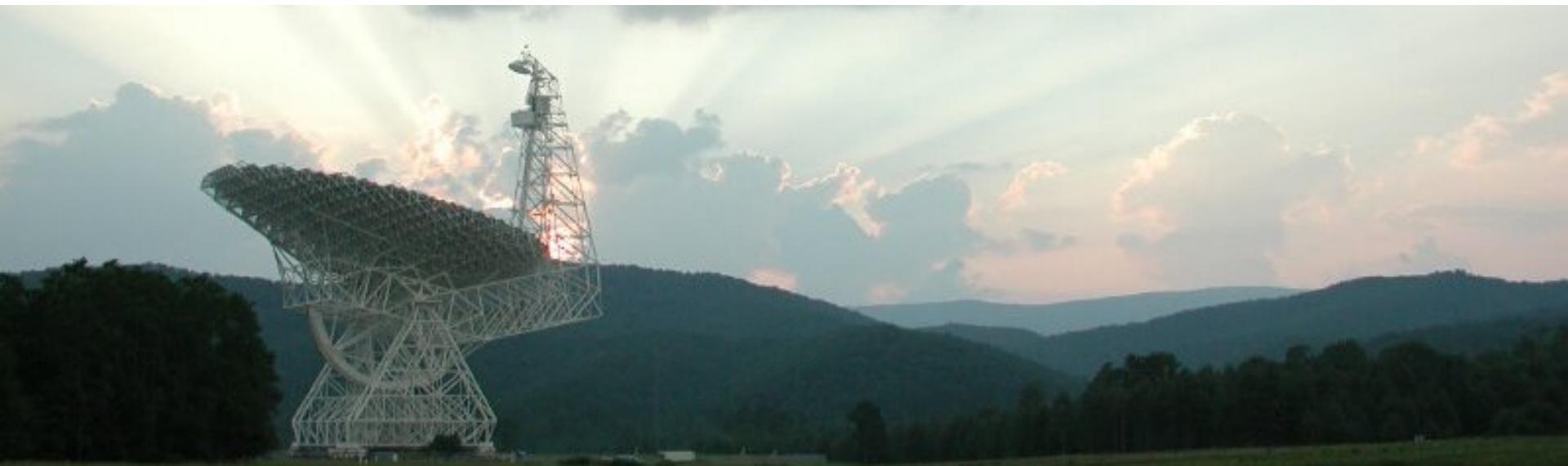


# Single Dish Observing Techniques and Calibration



David Frayer (NRAO)

{some slides taken from past  
presentations of Ron Maddalena and  
Karen O'Neil}



# What does the telescope measure:

$T_a$  = “antenna temperature”

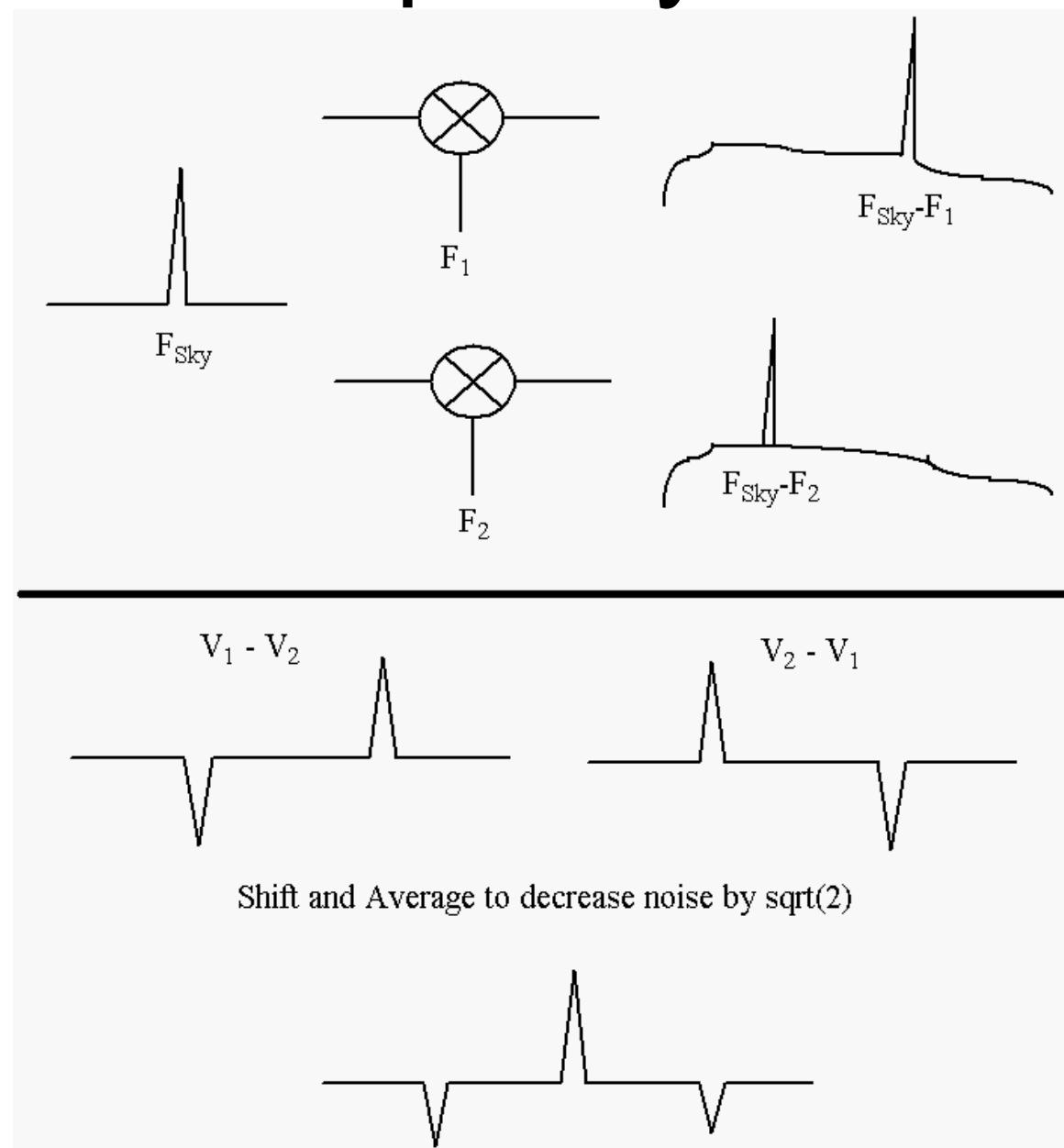
- $T_a(\text{total}) = T_{\text{source}} + \{T_{\text{rx}} + T_{\text{bg}} + T_{\text{atm}} + T_{\text{spill}}\}$
- Where  $\{....\}$  = other contributions
- Want  $T_{\text{source}}$ , so carry out ON – OFF
- $T_a(\text{ON}) = T_{\text{source}} + \{....\}$
- $T_a(\text{OFF}) = \{....\}$
- So  $T_a(\text{ON}) - T_a(\text{OFF}) = T_{\text{source}}$

→ Need to carry out ON-OFF observations and there are different observing techniques for measuring ON-OFF

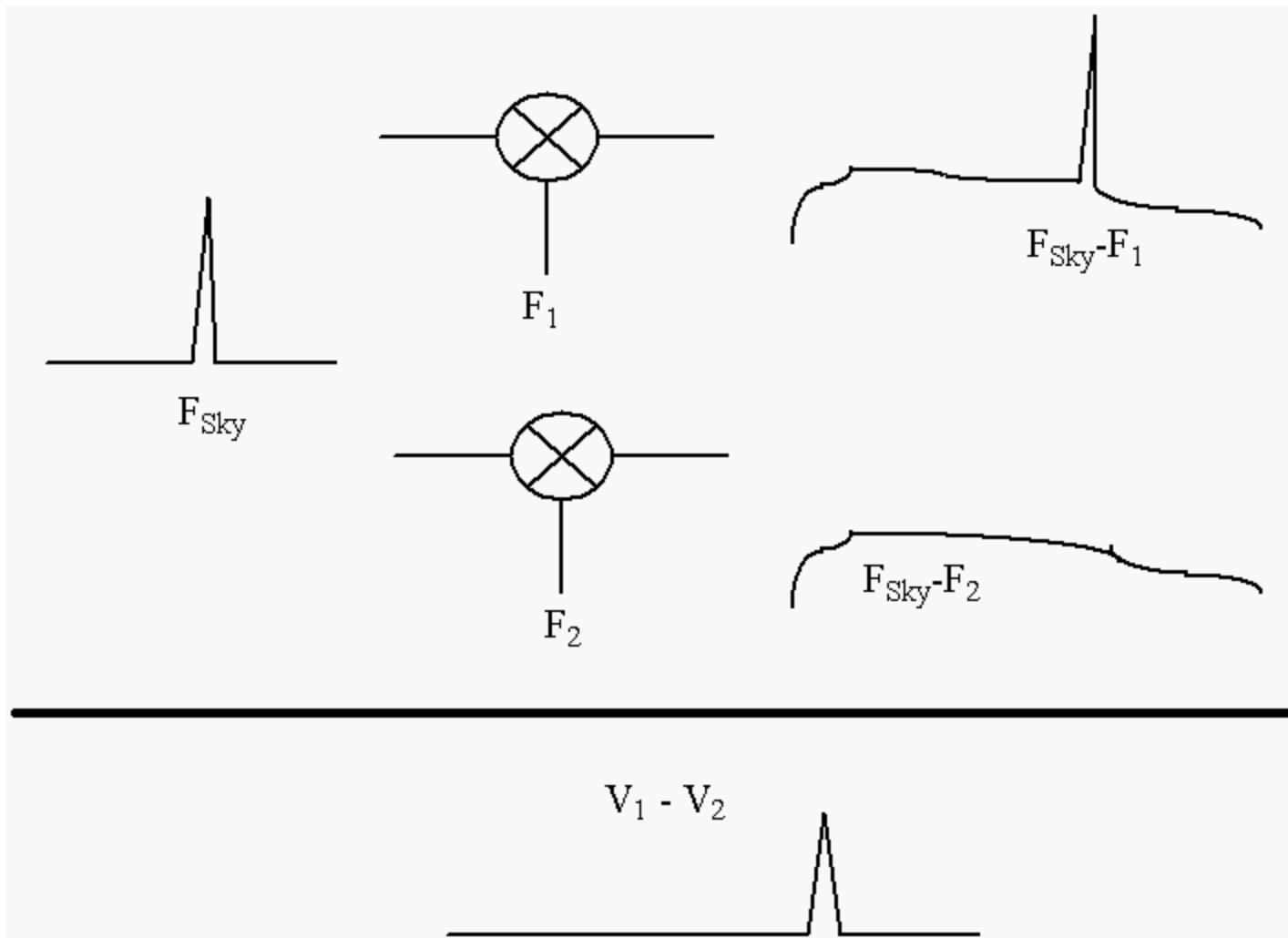
# Different Observing Modes to derive the reference data (OFF)

- Types of reference observations
  - Frequency Switching
    - In or Out-of-band
  - Position Switching
  - Beam Switching
    - Move Subreflector
    - Receiver beam-switch
  - Dual-Beam Nodding
    - Move telescope
    - Move Subreflector

# In-Band Frequency Switching



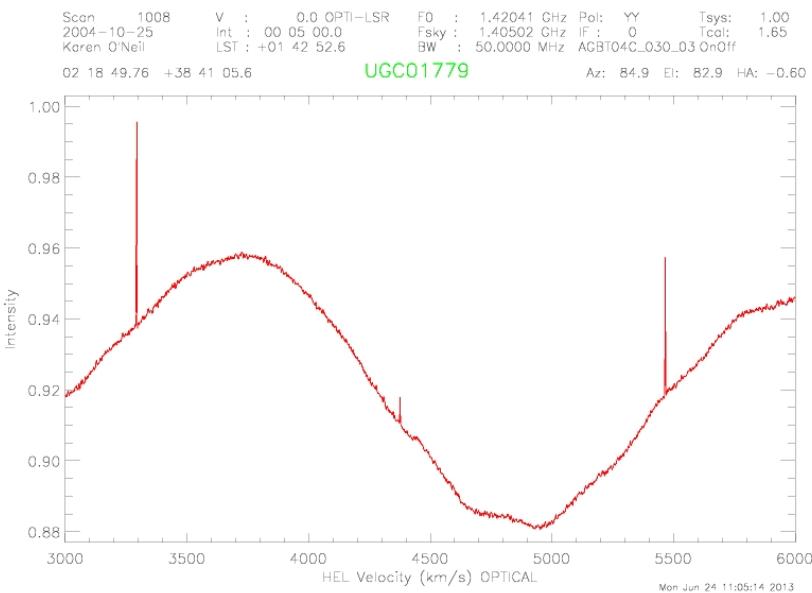
# Out-Of-Band Frequency Switching



# Position Switching

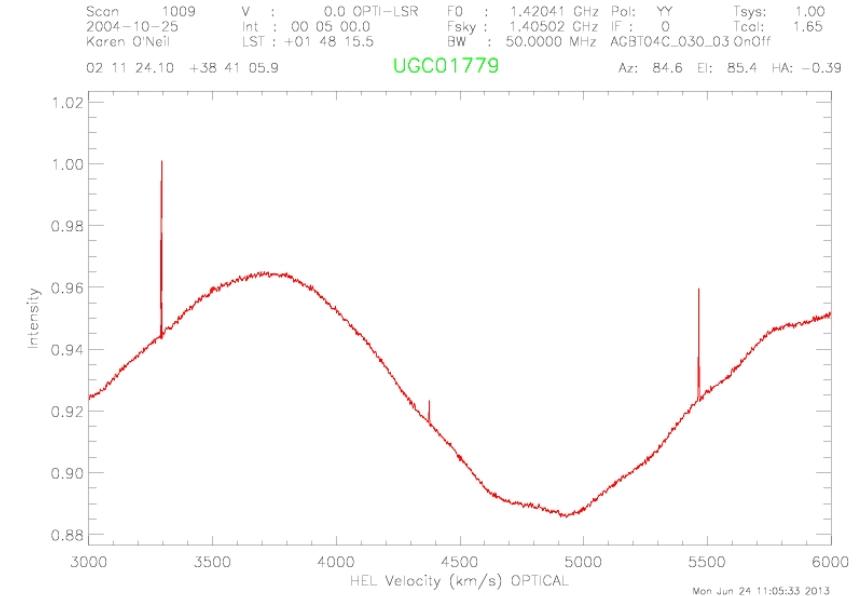
ON source

$$T_{\text{source}} + T_{\text{everything else}}$$



OFF source

$$T_{\text{everything else}}$$

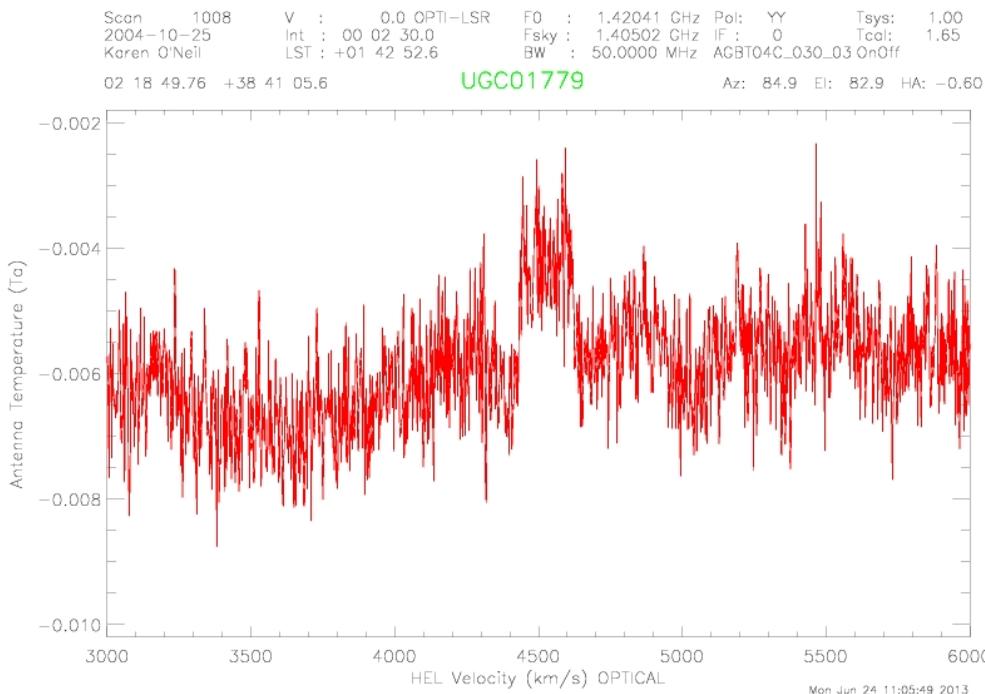


# Position Switching: ON-OFF on Sky

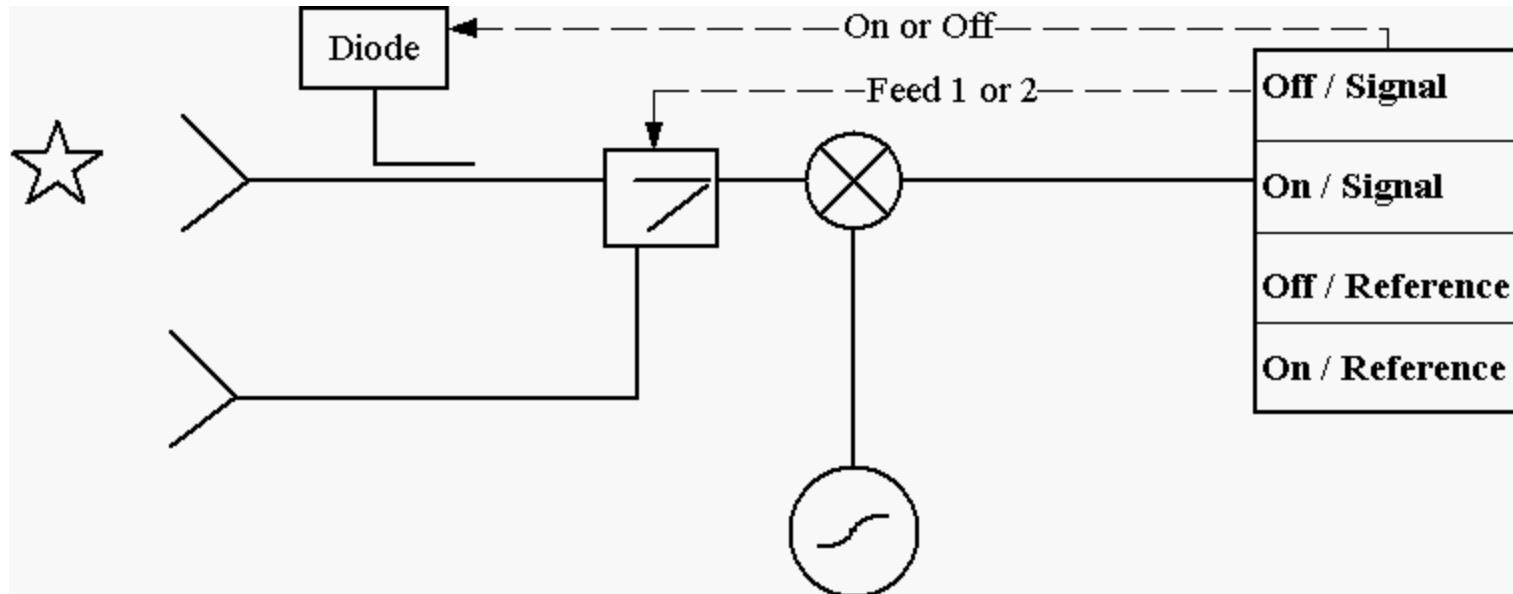
ON - OFF

$$(T_{\text{source}} + T_{\text{everything else}}) - (T_{\text{everything else}})$$

Arbitrary Counts

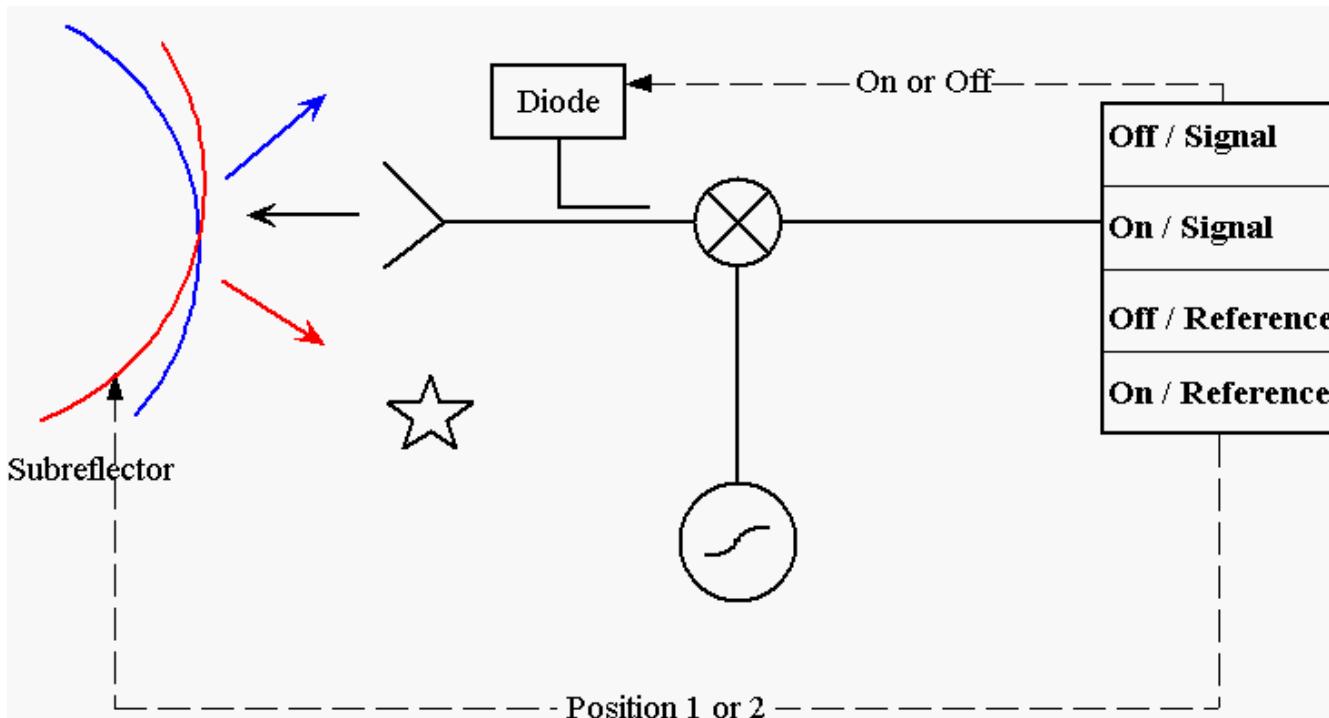


# Beam switching – Internal switching (“Dicke switching”)



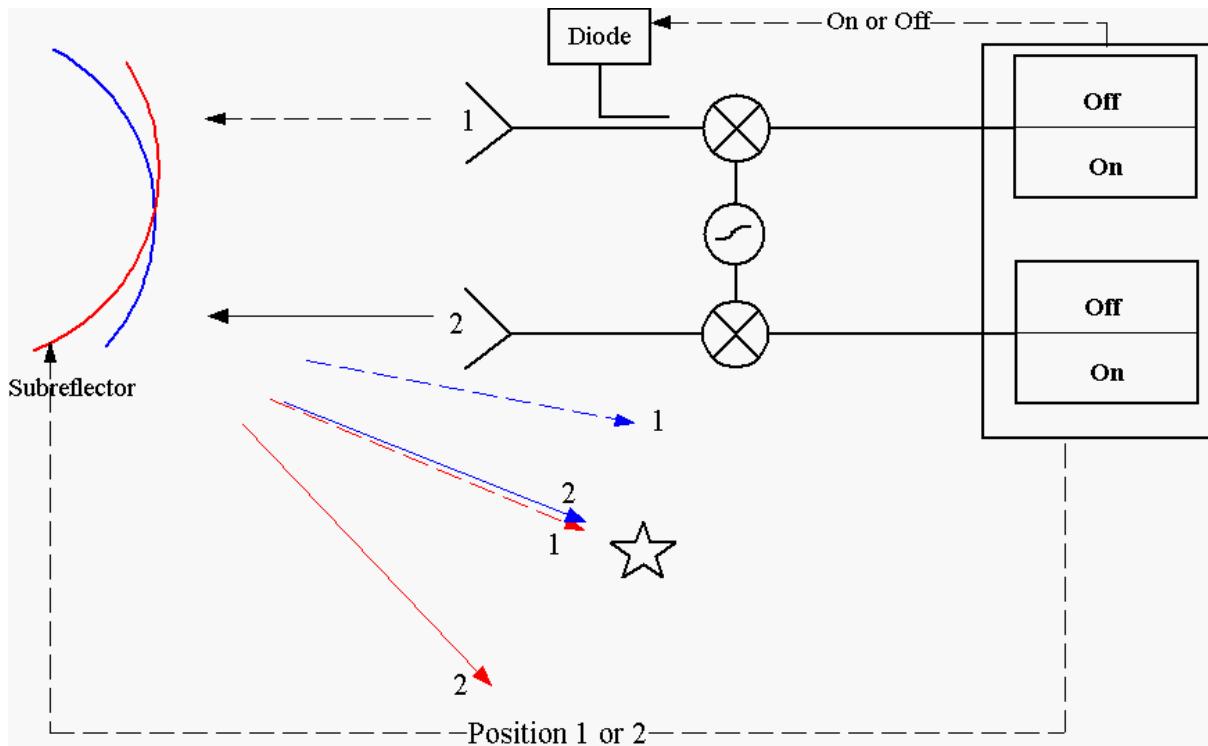
- Difference spectra eliminates any contributions to the bandpass from after the switch
- Residual will be the difference in bandpass shapes from all hardware in front of the switch.
- Low overhead but  $\frac{1}{2}$  time spent off source

# Beam Switching – Subreflector or tertiary mirror



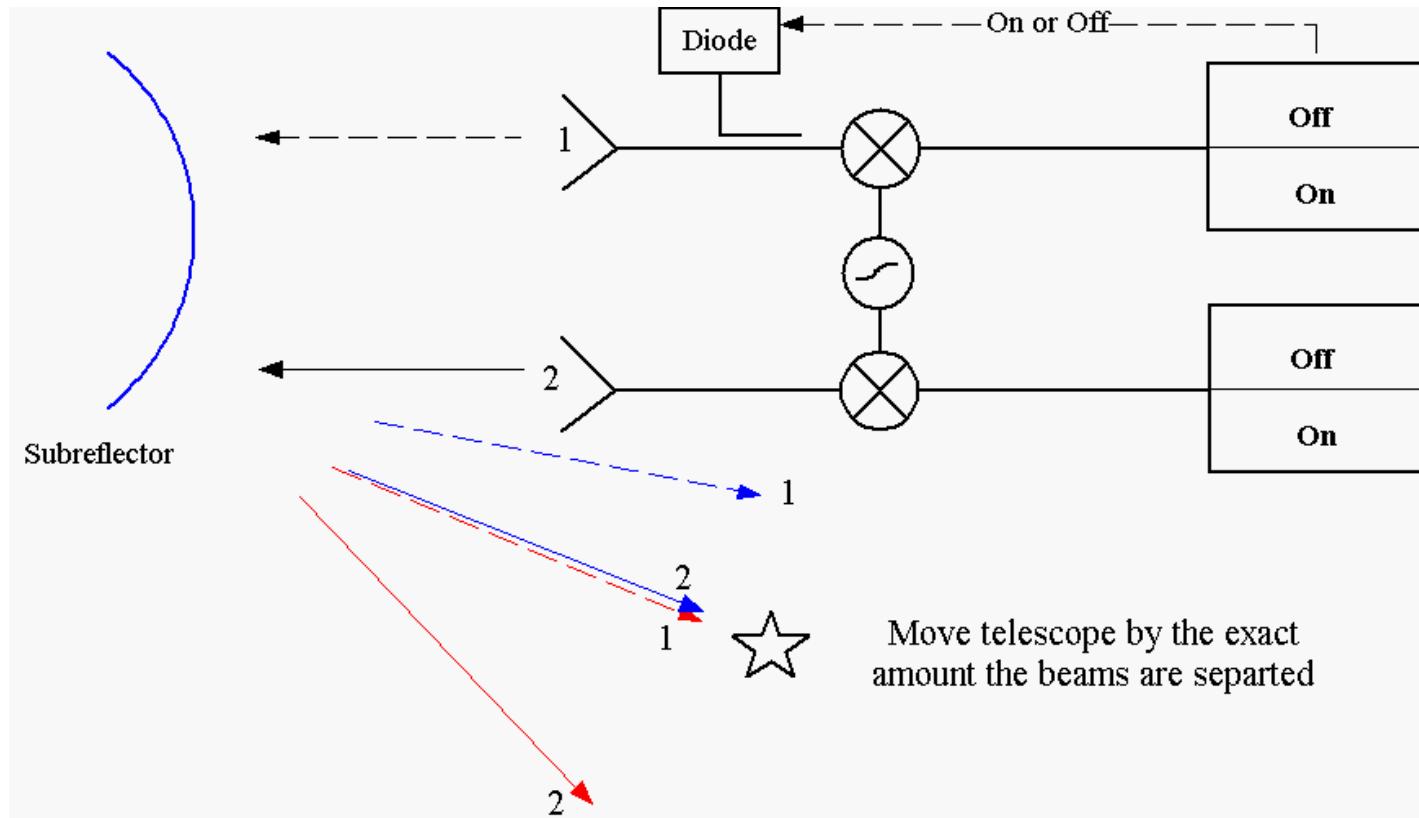
- Removes any ‘fast’ gain/bandpass changes
- Low overhead.  $\frac{1}{2}$  time spent off source

# Nodding with dual-beam receivers - Subreflector or tertiary mirror (SubBeamNod)



- Removes any ‘fast’ gain/bandpass changes
- Low overhead. All the time is spent on source

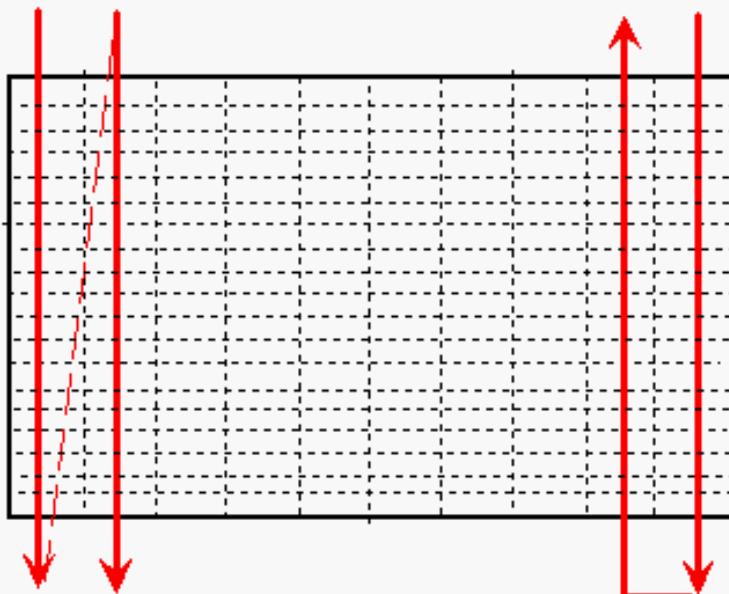
# Nodding with dual-beam receivers - Telescope motion (NOD)



- Removes any ‘fast’ gain/bandpass changes
- Overhead from moving the telescope. All the time is spent on source

# Mapping Techniques

1	7					1
2	8					2
3	9					3
4					9	4
5					8	5
6					7	6



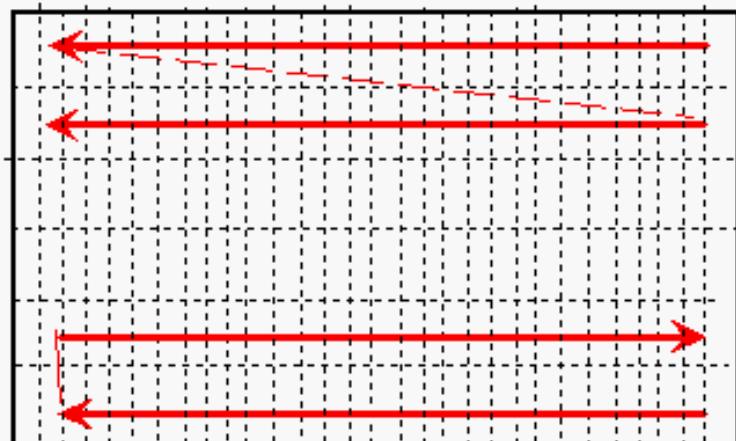
- **Point map**

- Sit, Move, Sit, Move, etc.

- **On-The-Fly Mapping**

- Slew a column or row while collecting data
- Move to next column/row
- Basket weave
- Should oversample ~3x Nyquist along direction of slew

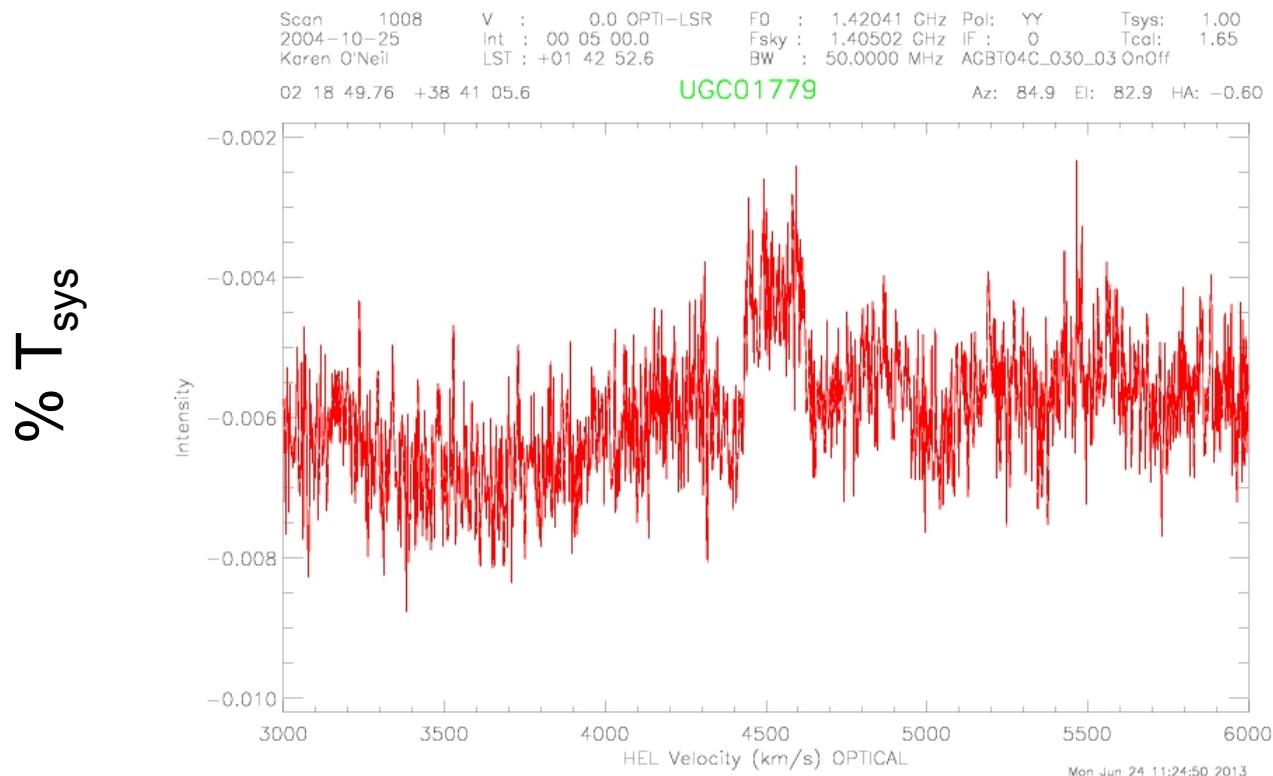
Reference/OFF from a “source-free” map position or separate “OFF” spectrum taken.



# Calibration of Data

(ON - OFF)/OFF

$$[(T_{\text{source}} + T_{\text{everything else}}) - (T_{\text{everything else}})] / T_{\text{everything else}}$$
$$= (\text{Source temperature}) / (\text{"System" temperature})$$



## Determining $T_a$

$$T_a = \frac{(ON - OFF)}{OFF} T_{\text{system}}$$

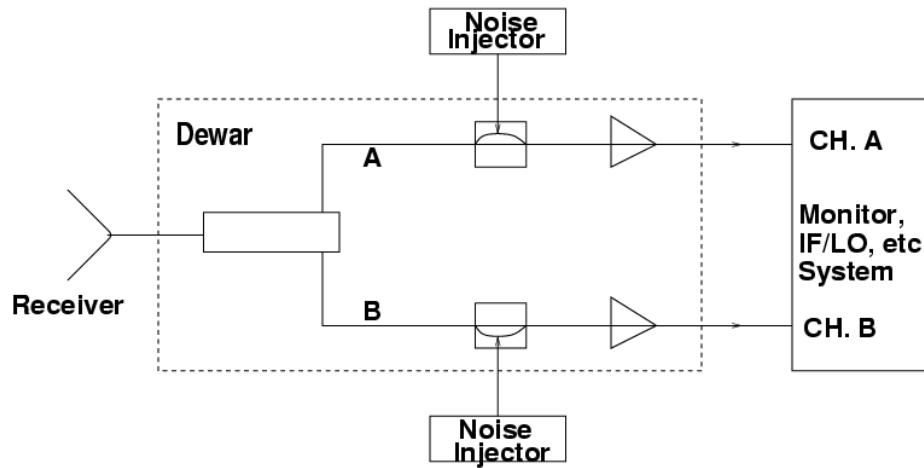
Blank Sky or other

From diodes, Hot/Cold loads, etc.

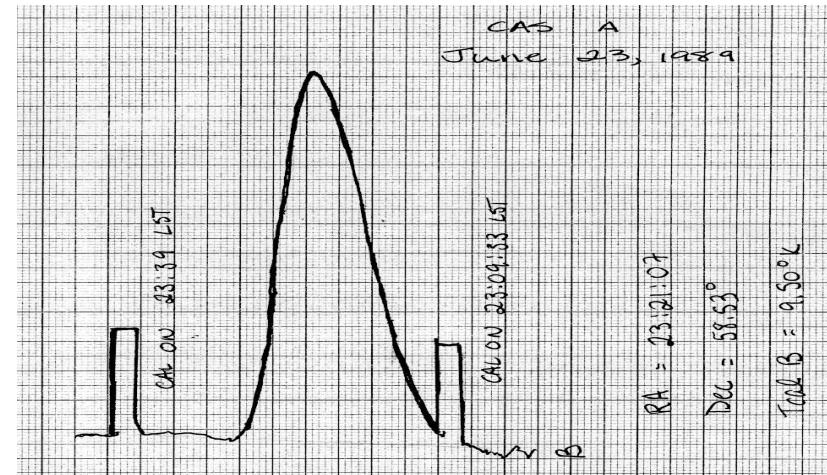
GBT definition of  $T_a$

# Determining $T_{sys}$

## Noise Diodes



All GBT receivers  
besides 4mm and  
Mustang use noise  
diodes.



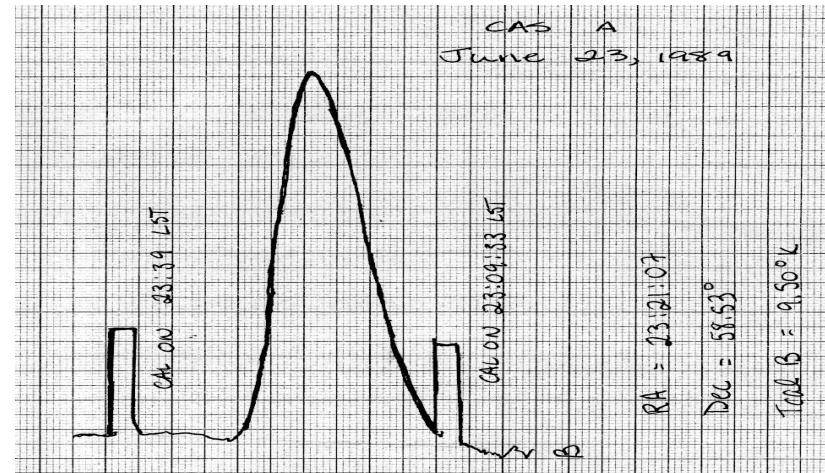
# Determining $T_{sys}$

## Noise Diodes

$$T_{sys} = T_{cal} * OFF / (ON - OFF)$$

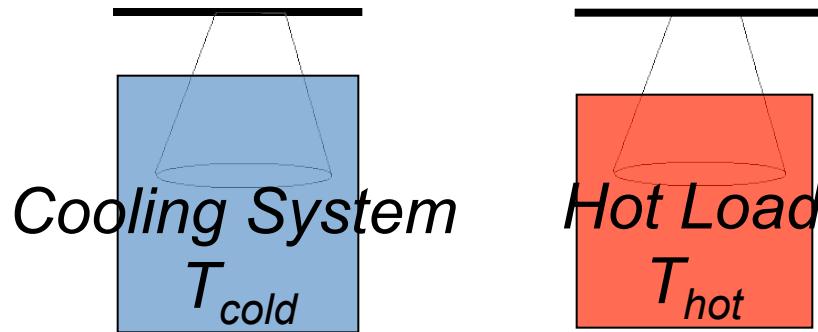
GBT: Flicker diode on/off

$$T_{sys} = T_{cal} * OFF / (ON - OFF) + T_{cal} / 2$$



# Determining $T_{sys}$

## Hot & Cold Loads



Gain:  $g = (T_{hot} - T_{cold}) / (V_{hot} - V_{cold})$  [K/Volts]

$$T_{sys} = g V_{off}$$

Example GBT 4mm Rx

# Absolute Calibration on known astronomical sources (point sources)

→ Corrects for any errors in the adopted Tdiode/gains measured in the lab and corrects for the telescope response

Observe and process source and known calibrator (3cX) source data in the same way, then the flux density of the source  $S(\text{source})$  is simply:

$$S(\text{source})/S(3\text{cX}) = T(\text{source})/T(3\text{cX}),$$

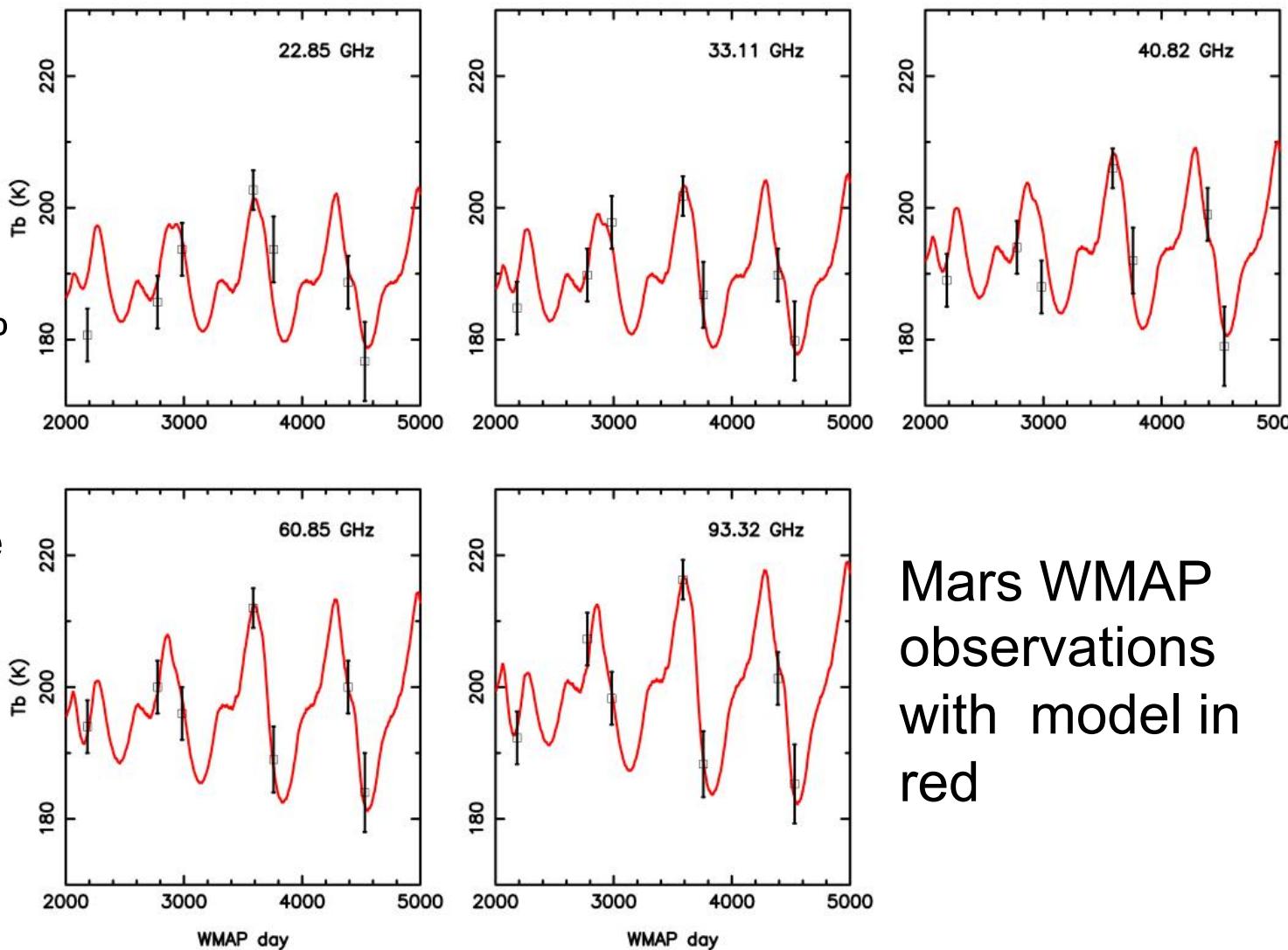
where  $S(3\text{cX})$  is known.

# Absolute Calibration tied to Mars via WMAP

VLA  
calibration  
(1-50 GHz):

- <20 GHz ~1% accurate
- 20-50 GHz: ~3% accurate

Perley &  
Butler 2013



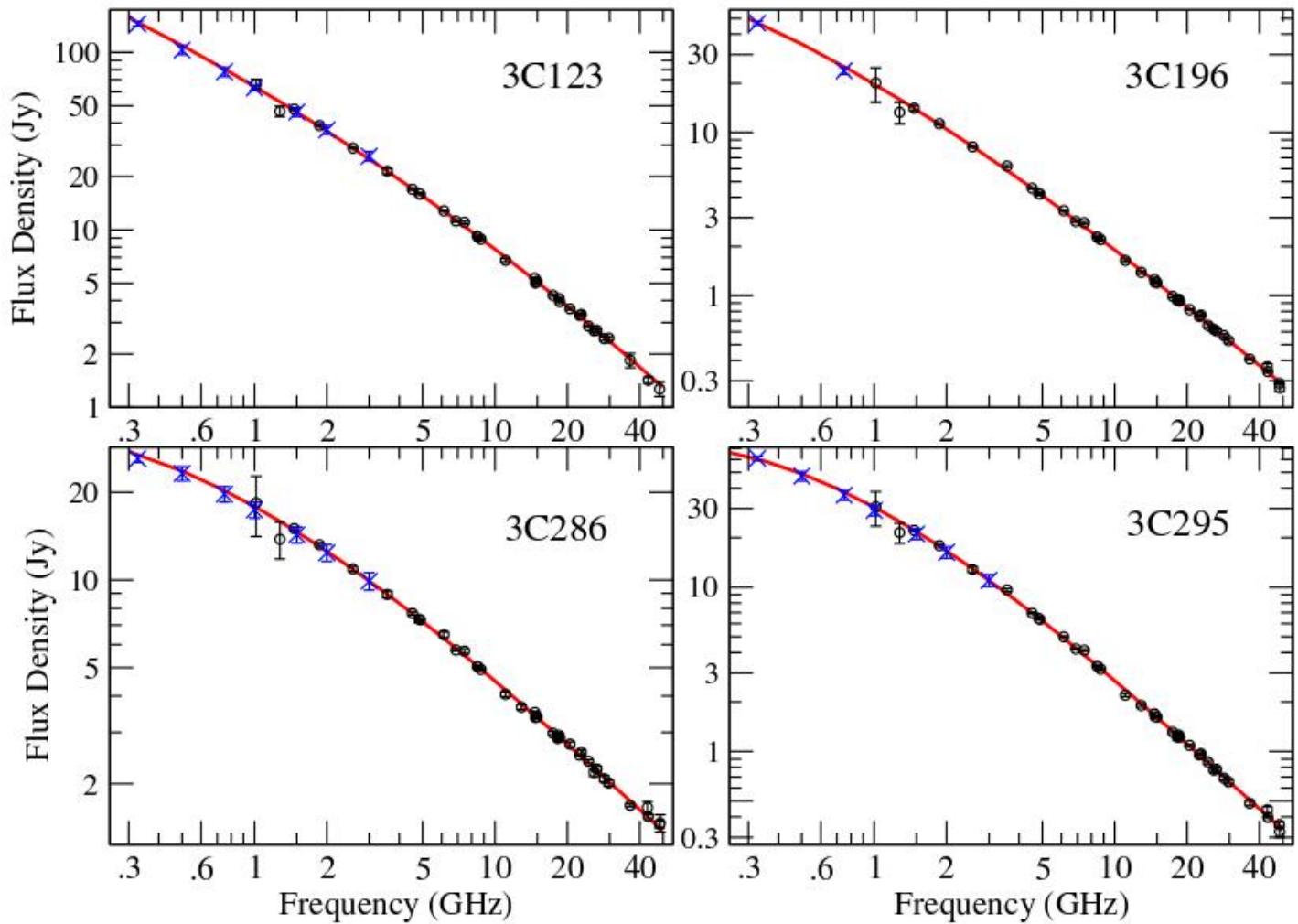
Mars WMAP  
observations  
with model in  
red

# VLA Stable Calibrators

GBT Calibration

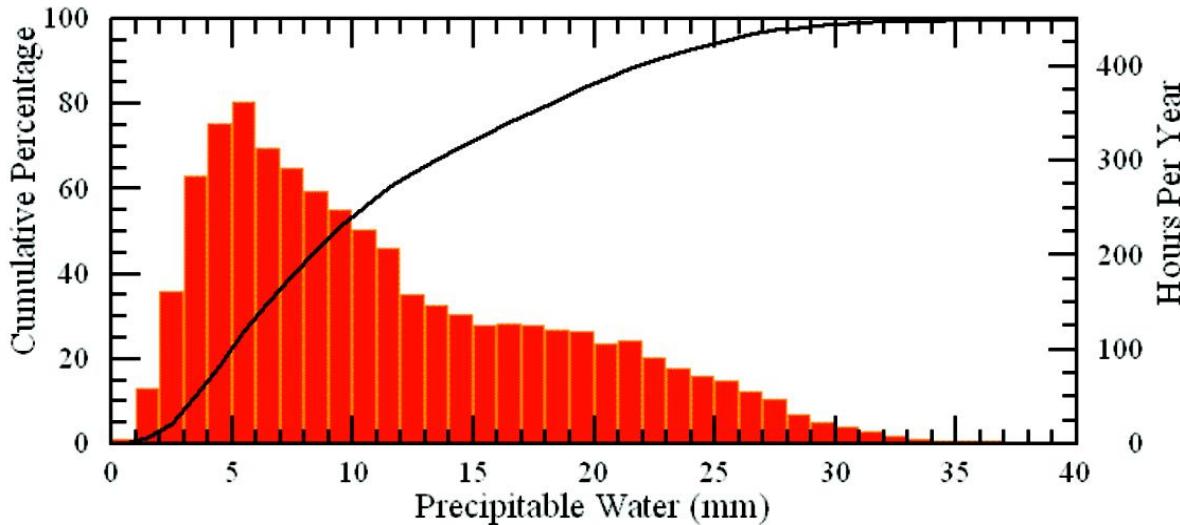
“Plan”:

Tie GBT to VLA calibration for 1-50 GHz, and we will use ALMA for 3mm absolute calibration (CARMA in the past)



# The atmosphere is important at high frequency (>10 GHz)

- Opacity
  - $T_{sys} = T_{rcvr} + T_{spill} + T_{bg} * \exp(-\tau * A) + T_{atm} * [\exp(-\tau * A) - 1]$
  - Air Mass  $A \sim 1/\sin(Elev)$  (for Elev > 15°)
- Stability
  - $T_{sys}$  can vary quickly with time
  - Worse when Tau is high

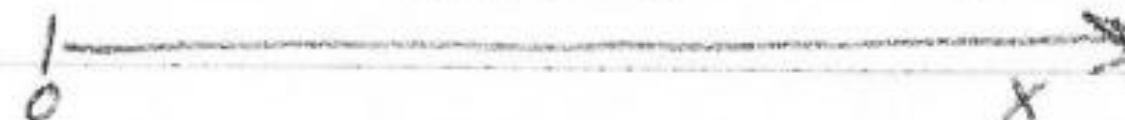


GBT site has many days with low water vapor per year (<10mm H<sub>2</sub>O are ok for 3mm, 50% of time)

# Radiative Transfer

Radiative Transfer

$$I_{\nu}(x) \rightarrow \overset{dx}{\text{---}} \rightarrow I_{\nu}(x+dx)$$



$$\frac{dI_{\nu}}{dx} = -K_{\nu} I_{\nu} + \epsilon_{\nu}$$

absorption      emission

# Thermal Equilibrium

$$\frac{dI_\nu}{dx} = -K_\nu I_\nu + \bar{\epsilon}_\nu$$

↑                      ↑  
absorption          emission

Thermal Equilibrium

$$\frac{dI_\nu}{dx} = 0, \quad I_\nu = \bar{B}_\nu(T) = \frac{\bar{\epsilon}_\nu}{K_\nu}$$

"Kirchhoff's law"

# Blackbody Equation

Blackbody eq:

$$B_{\nu}(T) = \frac{2h\nu^3}{c^2} \left[ \frac{1}{e^{h\nu/kT} - 1} \right]$$

$$B(T) = \int B_{\nu}(T) d\nu = \sigma T^4$$

"Stefan-Boltzmann"  
law

# Rayleigh-Jeans Approximation

Rayleigh-Jeans Approximation:  
 $h\nu \ll kT$

$$\Rightarrow B_{\nu}(T) \simeq \frac{2\nu^2}{c^2} kT = \frac{2kT}{\lambda^2} = I_{\nu} \text{ (Thermal-RS)}$$

Good for 3mm and longer wavelengths

# Temperature Scales

- $T_a = T_{sys} (\text{ON-OFF})/\text{OFF}$  (uncorrected antenna temperature)
- $T_a' = T_a \exp(\tau_o A)$
- $T_{mb} = T_a'/\eta_{mb}$  ( $\eta_{mb} \sim 1.3 \eta_a$ )
- $T_a^* = T_a'/\eta_I$  (mm-telescopes typically return  $T_a^*$ )
- $T_r^* = T_a' / (\eta_I \eta_{fss})$
- $T_a'/Sv = 2.84 \eta_a$  (for the GBT)

# Antenna Theorem

Power Received = Power Transmitted

$$P_{\text{rec}}(\theta, \phi) = \frac{1}{2} A_e P_n(\theta, \phi) S_v \Delta v \quad (\frac{1}{2} \text{ from single pol. })$$

$$S_v = \text{flux density [W/m}^2/\text{Hz}] = I_v \delta\Omega = (2kT/\lambda^2) \delta\Omega$$

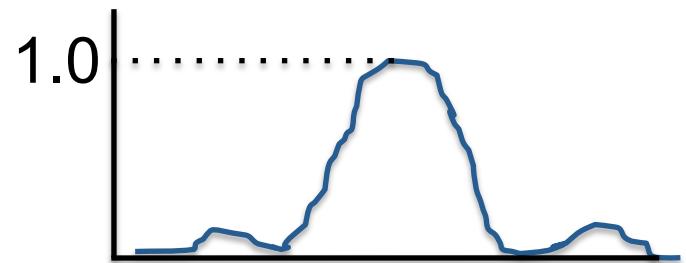
$$P_{\text{rec}} = \frac{1}{2} A_e \Delta v (2kT/\lambda^2) \iint P_n(\theta, \phi) d\Omega$$

$$P_{\text{trans}} = k T \Delta v$$

$$\frac{1}{2} A_e \Delta v (2kT/\lambda^2) \Omega_a = k T \Delta v$$

$$\rightarrow A_e \Omega_a = \lambda^2, \text{ where}$$

$$\text{Antenna Solid angle: } \Omega_a = \frac{\iint P_n(\theta, \varphi) d\Omega}{4\pi}$$



$P_n$  = antenna power pattern normalized to the peak;  
 $P_n(0,0)=1.0$

# Point-Source Calibration: Flux Density vs Antenna Temp

$$P_{\text{rec}} = \frac{1}{2} A_e S_v \Delta v = k T_a' \Delta v$$

$$A_e = \eta_a (\pi/4) D^2$$

$$\rightarrow S_v = 3520 T_a' / (\eta_a [D/m]^2)$$

i.e.,  $T_a' / S_v = 2.84 \eta_a$  for the GBT ( $\eta_a=0.71$  at low  $v$ )

Used for point-source calibration:

- Measure  $T_a$
- Correct for atmosphere  $\rightarrow T_a'$
- Know  $S_v$
- Derive  $\eta_a$

# Telescope “Gain”

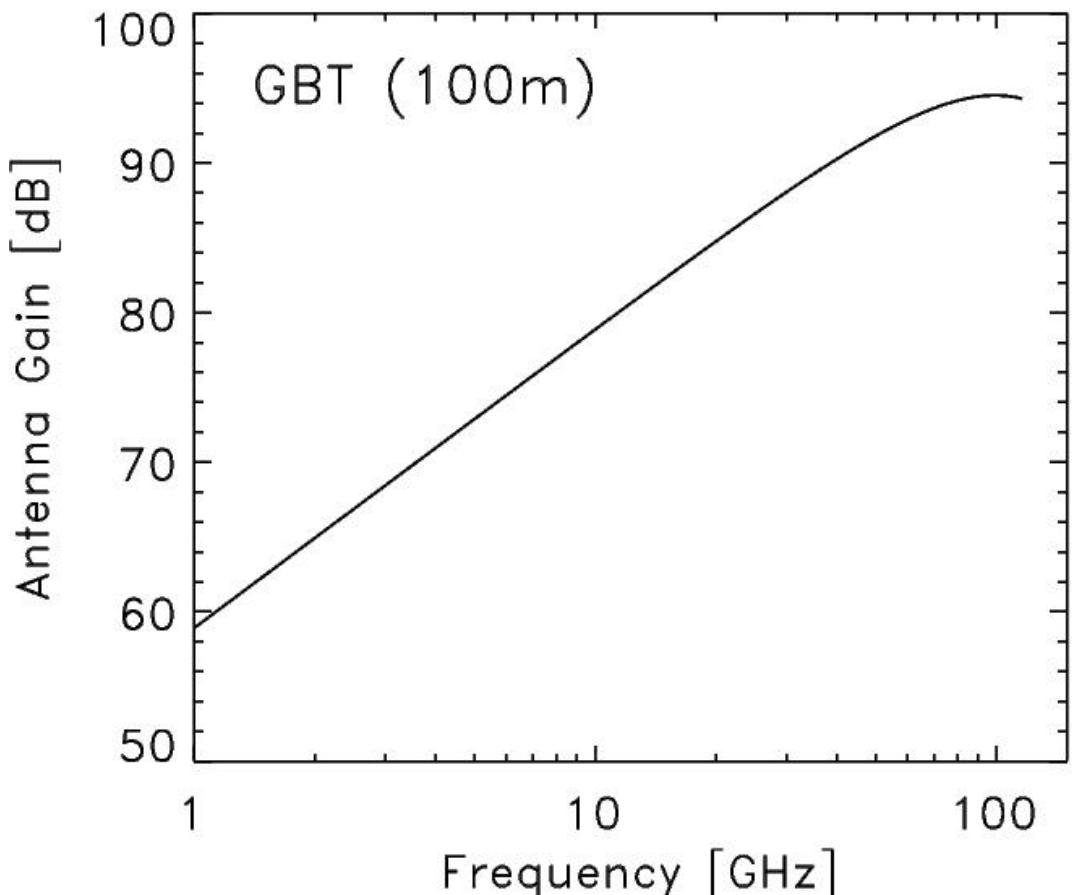
Astronomers:

$$\text{Gain} = T_a' / S_v \text{ in units of K/Jy}$$

Engineers: Gain as given by the Antenna Theorem

$$\rightarrow G = 4\pi A_e / \lambda^2 \text{ in units of dB}$$

(antenna viewed as a big amplifier)



$$T_R^* = T_R^* \eta_r \left[ \frac{\iint_{\Omega_d} P_n d\Omega}{4\pi} / \frac{\iint P_n d\Omega}{4\pi} \right]$$

↑

Kutner + Ulrich 1981

12m, "US" definitions

$\eta_r \approx 1$  radiative efficiency

Fraction of power  
not lost in spillover  
+ scattering

$$\eta_{fss} \equiv \frac{\iint_{\Omega_d} P_n d\Omega}{4\pi} / \frac{\iint P_n d\Omega}{2\pi}$$

"Forward" scattering  
+ spillover eff.

$$\eta_{rss} \equiv \frac{\iint_{\Omega_d} P_n d\Omega}{2\pi} / \frac{\iint P_n d\Omega}{4\pi}$$

"Rearward" scattering  
+ spillover eff.

Forward hemisphere

$$\Rightarrow \bar{T}_A' = \bar{T}_R * \eta_r \eta_{fss} \eta_{rss} = \bar{T}_R * \eta_{fss} \eta_e$$

$$\eta_e = \eta_r \eta_{rss}$$

$\eta_e = F_{\text{eff}}$  "forward efficiency" IRAM/Class

$\eta_e$  = fraction of power on sky in forward direction

$(1 - \eta_e)$  = fraction of power that "sees" the ground

$T_A^*$  "Chopper-wheel" Calibration gives  $T_a^*$

$$\bar{T}_a^* = \bar{T}_a' / \eta_l = \eta_{fss} \cdot \bar{T}_R^*$$

IRAM:  $\bar{T}_{mb} = \frac{F_{eff}}{B_{eff}} \bar{T}_a^*$

$$\bar{T}_{mb} = \frac{\bar{T}_a'}{\eta_{mb}} ; \bar{T}_a' = \bar{T}_a^* \eta_l$$

$$\bar{T}_{mb} = \bar{T}_a^* \underbrace{\eta_l}_{\eta_{mb}} \quad \eta_l = F_{eff}$$

$$B_{eff} = \eta_{mb}$$

# Calibration Two Loads vs One Load

Two loads "Direct" calibration, e.g. 4m Rx

$$g = \frac{(T_{amb} - T_{cold})}{(V_{amb} - V_{cold})}$$

$$\bar{T}_A = T_{sys} \left( \frac{ON - OFF}{OFF} \right) = g (ON - OFF)$$

$$\bar{T}_A' = \bar{T}_A e^{\bar{\tau}_0 A} \quad \text{need } \underline{\underline{\tau}}$$

# Calibration with One Load, $T_A^*$

One load

chopper calibration:

With a chopper wheel/ vane and a simple temperature sensor, one can calibrate to the approximate  $T_A^*$  scale without any knowledge of the sky.

$$T_A^* = T_{\text{cal}} \left[ \frac{V_{on} - V_{off}}{V_{amb} - V_{off}} \right]$$

need  $\overline{T}_{\text{ATM}}$

$$T_{\text{cal}} = \left[ \frac{T_{\text{amb}} - T_A^{\text{sky}}}{\eta_l} \right] e^{T_{\text{cal}}}$$

; yuck, but:

Algebra  $\Rightarrow$

$$T_{\text{cal}} = T_{\text{ATM}} + (T_{\text{amb}} - T_{\text{ATM}}) e^{T_{\text{cal}}}$$

what Argus will use.

↑  
Model

eq. A9 of KU91  
rewritten

assumes  $T_{\text{spill}} \approx T_{\text{amb}}$  and ignoring  $T_{\text{bg}}$

Some have assumed  $T_{\text{cal}} = T_{\text{amb}} = T_{\text{ATM}}$  with  $T_A^*$  calibration

Tsys for  $T_A^*$  scale different than Tsys for  $T_A$

$$\bar{T}_A^* = \bar{T}_{\text{sys}}^* \left( \frac{\text{ON-OFF}}{\text{OFF}} \right)$$

where  $\bar{T}_{\text{sys}}^* = \bar{T}_{\text{sys}} \cdot \frac{e^{\zeta_0 A}}{\eta_l}$  includes ATM

i.e

$$\bar{T}_A^* = \bar{T}_A \cdot \frac{e^{\zeta_0 A}}{\eta_l}$$

# Definitions of $\Omega$ 's

$$\Omega_A \equiv \iint_{4\pi} P_n(\theta, \phi) d\Omega$$

$$\eta_{mb} = \frac{\Omega_{mb}}{\Omega_A}$$

$$\Omega_{mb} \equiv \iint_{MB} P_n(\theta, \phi) d\Omega$$

$$\eta_{rss} = \frac{\Omega_F}{\Omega_A} \simeq \eta_F = F_{eff}$$

$$\Omega_d \equiv \iint_{2d} P_n(\theta, \phi) d\Omega$$

$$\eta_{fss} = \frac{\Omega_d}{\Omega_F}$$

$$\Omega_F \equiv \iint_{2\pi} P_n(\theta, \phi) d\Omega$$

$$\Omega_{source} \equiv \iint_{source} P_n(\theta, \phi) d\Omega$$

Note:

$$\Omega_d = \Omega_{mb}$$

(different authors/  
conventions)

# Extended Sources: $T_{mb}$ vs $T_{source}$

$$T_A' = \frac{1}{\Omega_A} \iint P_n(\theta, \phi) T_S(\theta, \phi) d\Omega$$

$\uparrow$   
Source

compute using  $T_{mb}$

$$\frac{\Omega_{mb}}{\Omega_A} = \eta_{mb} \quad ; \quad T_A' = \eta_{mb} T_{mb}$$

$$T_{mb} = \overline{T_A'} = \frac{1}{\Omega_{MB}} \iint P_n(\theta, \phi) T_S(\theta, \phi) d\Omega$$

case 1: If  $T_s$  is uniform,  $\theta_s < \theta_{mb} \Rightarrow \Omega_s \ll \Omega_{mb}$   
 $\Rightarrow P_n \approx 1$  over source:

$$T_{mb} = \frac{1}{\Omega_{mb}} T_s \iint_{\text{source}} P_n d\Omega$$

$$= T_s \frac{\Omega_s}{\Omega_{mb}} = \left( \frac{\theta_s}{\theta_{mb}} \right)^2$$

$$T_{mb} = T_s \left( \frac{\theta_s}{\theta_{mb}} \right)^2$$

Small source  $\theta_s < \theta_{mb}$



"Filling factor"  
 "Beam definition"

$$T_{mb} = \frac{1}{Q_{mb}} \iint P_n \cdot T_S d\Omega$$

Case 2:  $\theta_S = \theta_{mb}$ ,  $T_S$  uniform

$$T_{mb} = T_S \frac{\iint P_n d\Omega}{Q_{mb}}$$

$$\boxed{T_{mb} \equiv T_S \quad \theta_S = \theta_{mb}}$$

How about  $\theta_S > \theta_{mb}$ :

$$\Rightarrow \boxed{T_{mb} \approx T_S \quad \theta_S > \theta_{mb}}$$

Note: For beams with significant side-lobes the measured  $T_a'$  increases with source size outside the main-beam and the relationship between  $T_{mb}$  and  $T_S$  depends on the details of how  $Q_{mb}$  is derived and the coupling of the source to the main beam.

case 3

# Gaussian Source

More general, assume  $T_s$  Gaussian and beam Gaussian

$$\Omega_{\text{Gaussian}} = 1.133 \theta_{\text{FWHM}}^2$$

$$T_{mb} = T_s \left[ \frac{\theta_s^2}{\theta_{mb}^2 + \theta_s^2} \right]$$

$$\theta_{mb} \gg \theta_s \quad T_{mb} = T_s \left( \frac{\theta_s}{\theta_{mb}} \right)^2 \checkmark$$

$$\theta_s \gg \theta_{mb} \quad T_{mb} = T_s \quad \checkmark$$

# Concluding Remarks

- To observe weak signals, one needs to measure ON-OFF
- Several different observing techniques can be used to give ON-OFF (freq-switched, position switched)
- At cm wavelengths, we use noise diodes to calibrate the data, while at mm wavelengths ambient/cold loads are used
- At low-freq, the Ta scale is used, while at high freq, one must correct for atmosphere (Ta')
- The Ta\* scale is typically used with one ambient load at mm wavelengths (chopper technique)
- Point sources are typically calibrated to flux density [Sv] Jy units, while extended sources are typically calibrated to the  $T_{mb}$  [K] temperature scale