

If Accretion Controls Filament Evolution...

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(4) its signature could still be hidden.

(1) Accretion and Fragmentation

Fragmentation timescale Nagasawa 87, Tomisaka 95

$$\tau_f = \frac{3}{\sqrt{4\pi G \rho_c}} = 5.24 \times 10^5 \left(\frac{n_c}{10^4 \text{ cm}^{-3}} \right)^{-1/2} \text{ yrs}$$

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Accretion timescale FH+ 09, Palmeirim+ 13

$$\tau_a = 8.31 \times 10^5 \left(\frac{T}{10 \text{ K}} \right)^{1/2} \left(\frac{n_{ext}}{100 \text{ cm}^{-3}} \right)^{-1} \left(\frac{R}{\text{pc}} \right)^{-1} \left(\ln \frac{R_{ref}}{R} \right)^{-1/2} \text{ yrs.}$$

based on steady-state accretion $v_R = 2 \left(Gm \ln \frac{R_{ref}}{R} \right)^{1/2}$

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based on steady-state accretion $v_R = 2 \left(\dot{M} \ln \frac{R_{ref}}{R} \right)^{1/2}$

Depends on line mass, filament radius.

(1) Accretion and Fragmentation

Line mass evolution:

$$\frac{dm}{dt} = 2\pi R \rho_{ext} v_R = 4\pi \rho_{ext} R \left(Gm(t) \ln \left(\frac{R_{ref}}{R} \right) \right)^{1/2}$$

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Setting the radius:

$$R_f = R_0 \left(\left(\frac{\rho_c}{\rho_{ext}} \right)^{2/p} - 1 \right)^{1/2} \quad \text{arbitrary } p \quad \rho(R) = \rho_c \left(1 + \left(\frac{R}{R_0} \right)^2 \right)^{-p/2}$$

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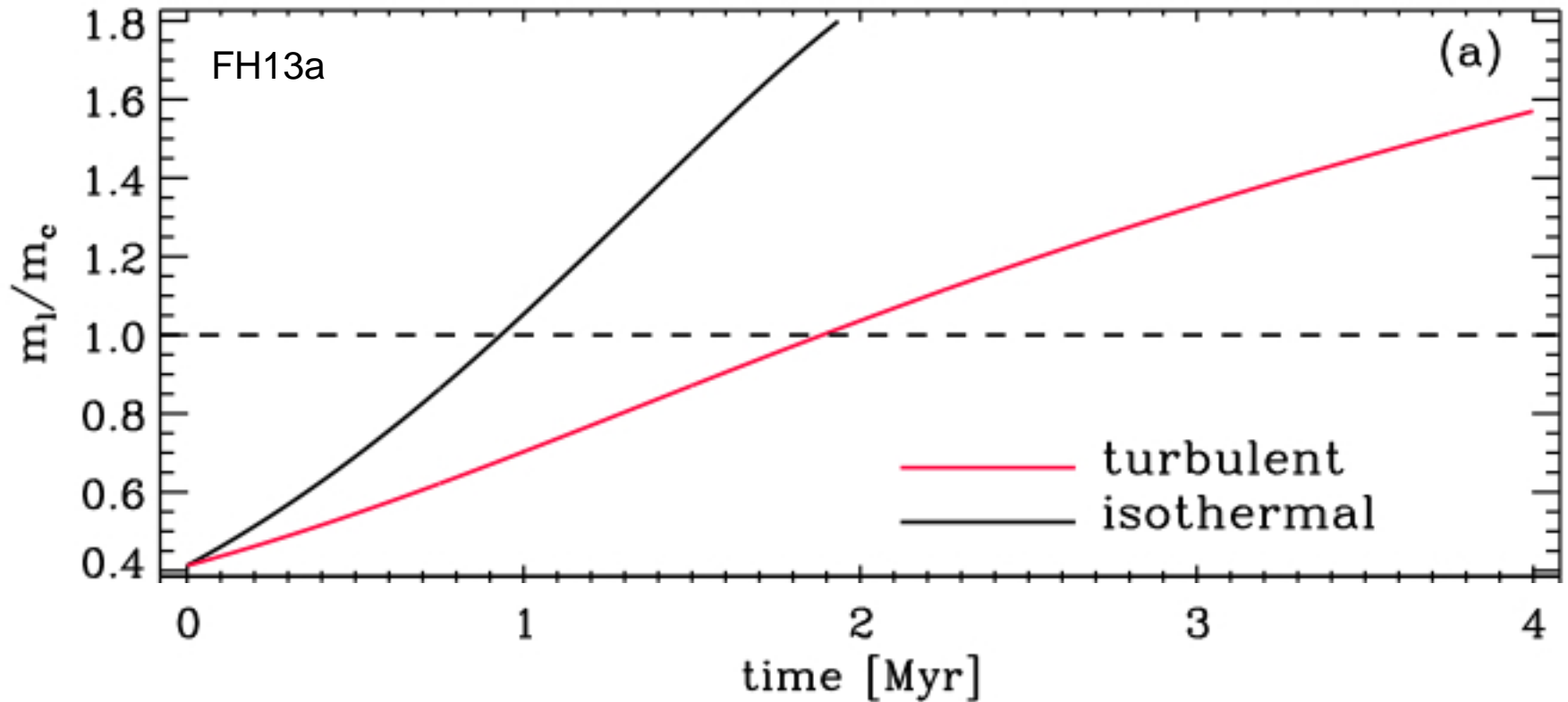
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$$R_f = \sigma^2 \left(\frac{2f(1-f)}{\pi G p_{ext}} \right)^{1/2} \quad \text{externally pressurized, isothermal}$$

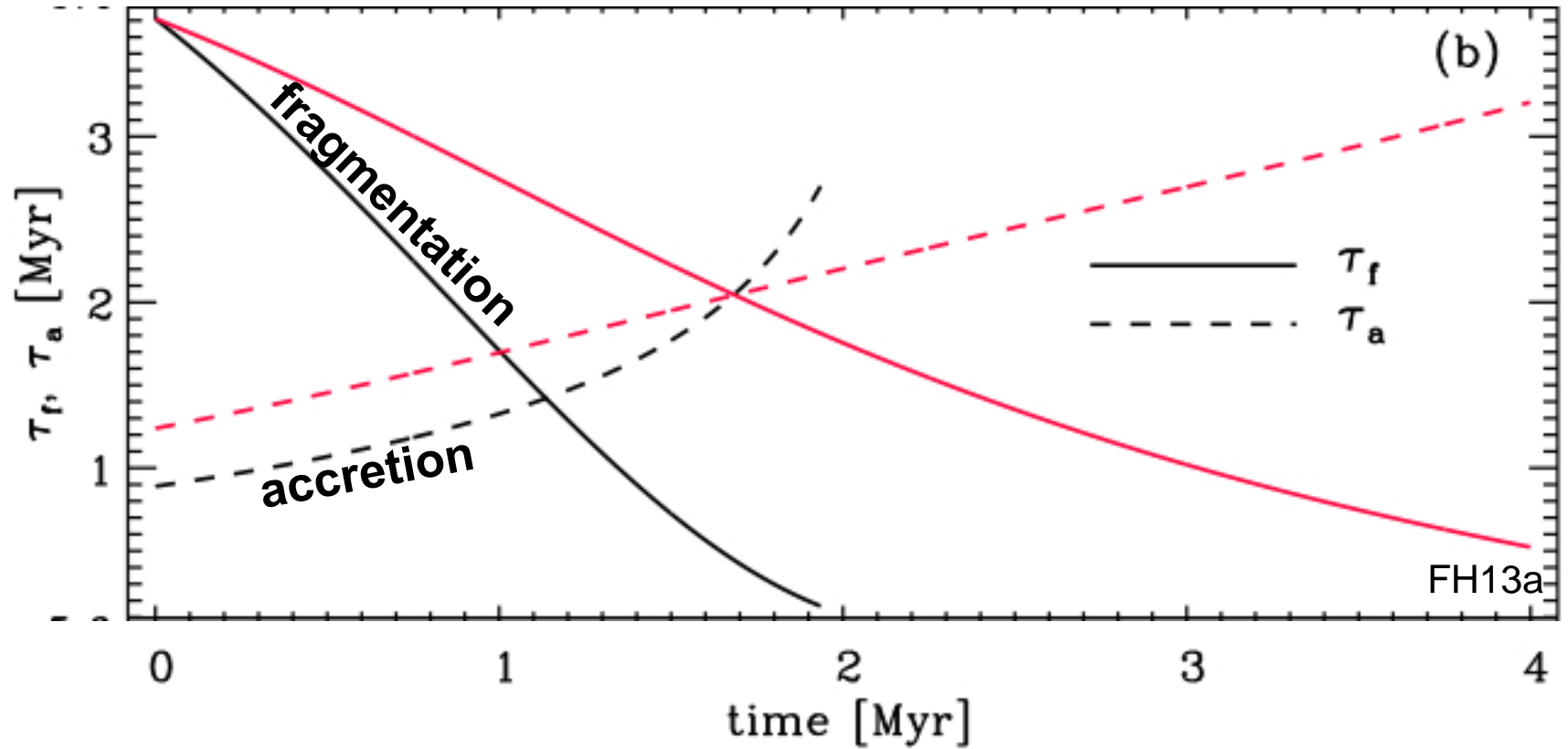
Fischera & Martin 12a,b

(1) Accretion and Fragmentation

line mass $m(t)$ for $p=1.5$



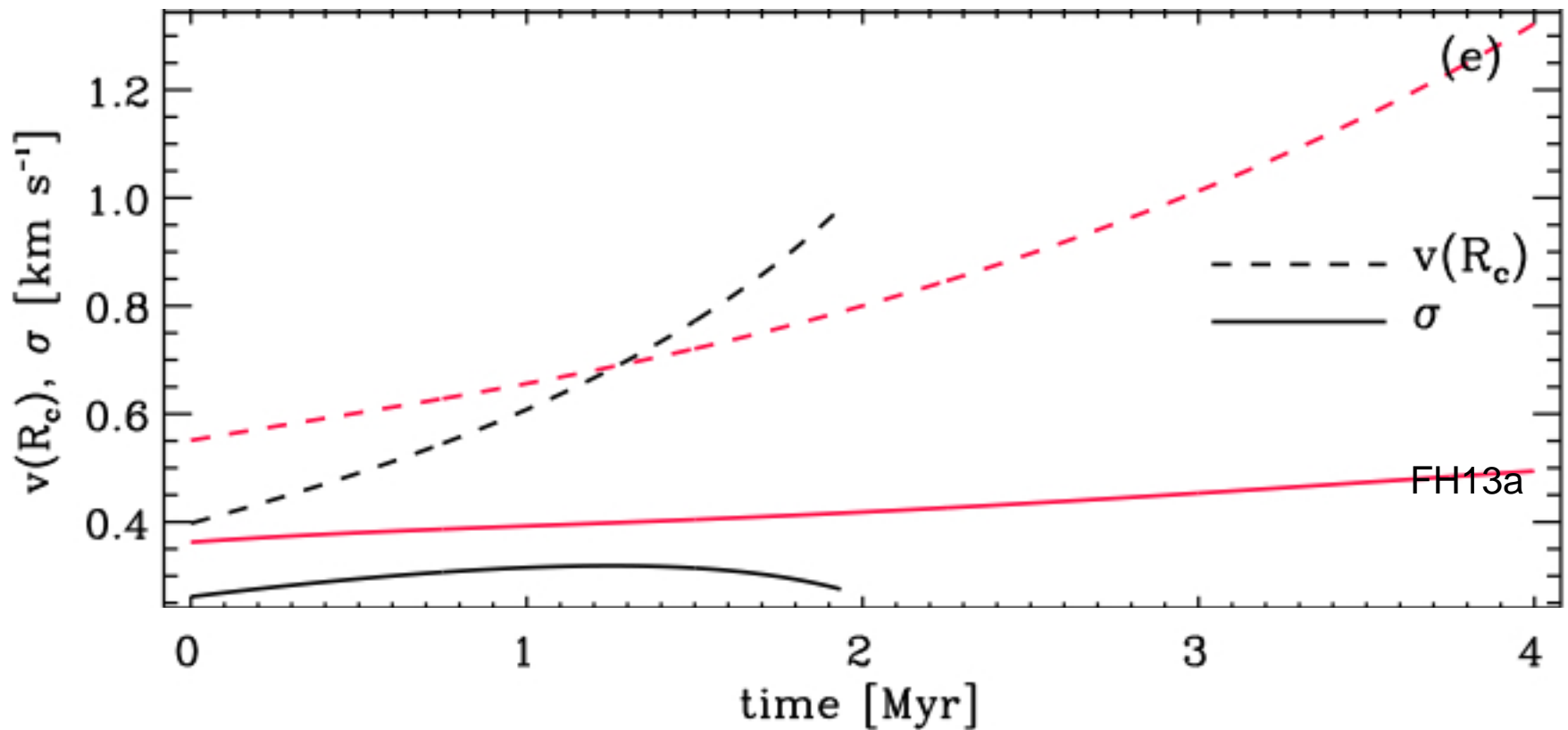
(1) Accretion and Fragmentation



(2) Accretion-driven Turbulence Klessen & Hennebelle 10

motivated by Vazquez-Semadeni+ 07, FH+ 08, Field+ 08...

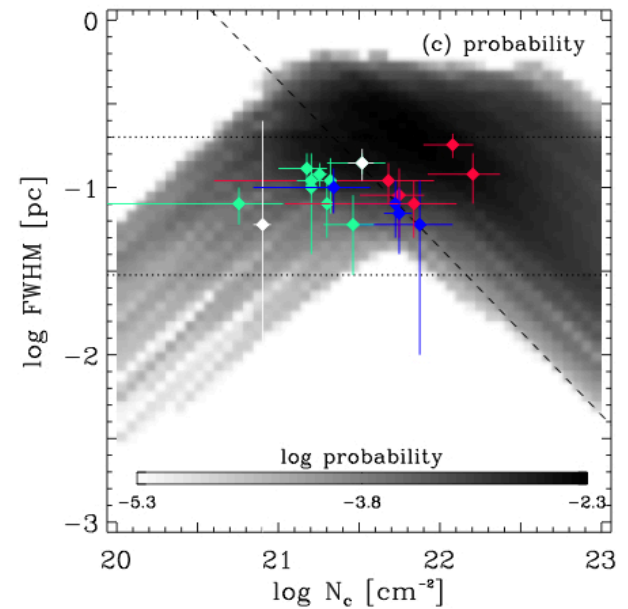
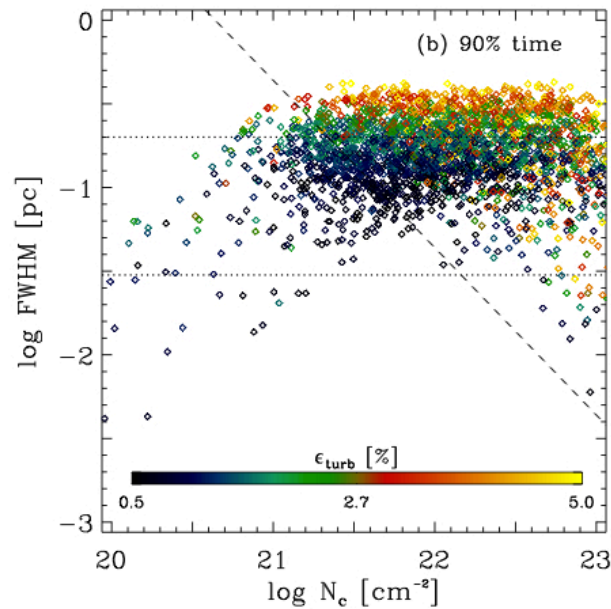
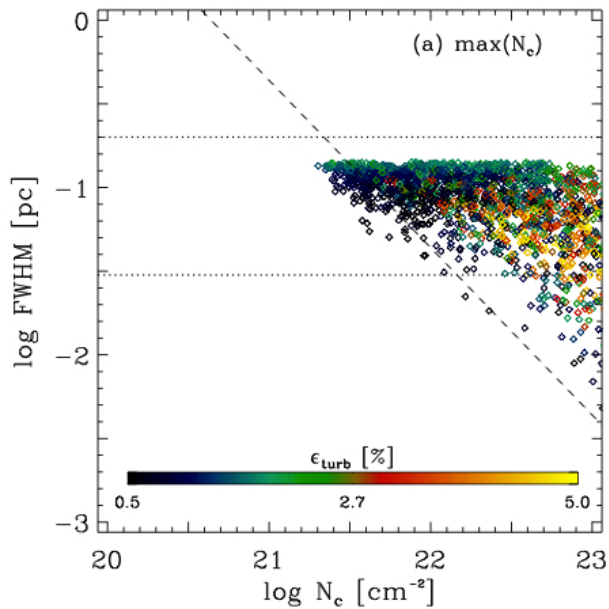
$$\sigma = \left(2\epsilon R_f v^2 (R_f) \frac{dm/dt}{m(t)} \right)^{1/3}$$



(2) Accretion-driven Turbulence

FWHM(N_c) relations (version 1: $p=1.5$)

[Arzoumanian+ 11]

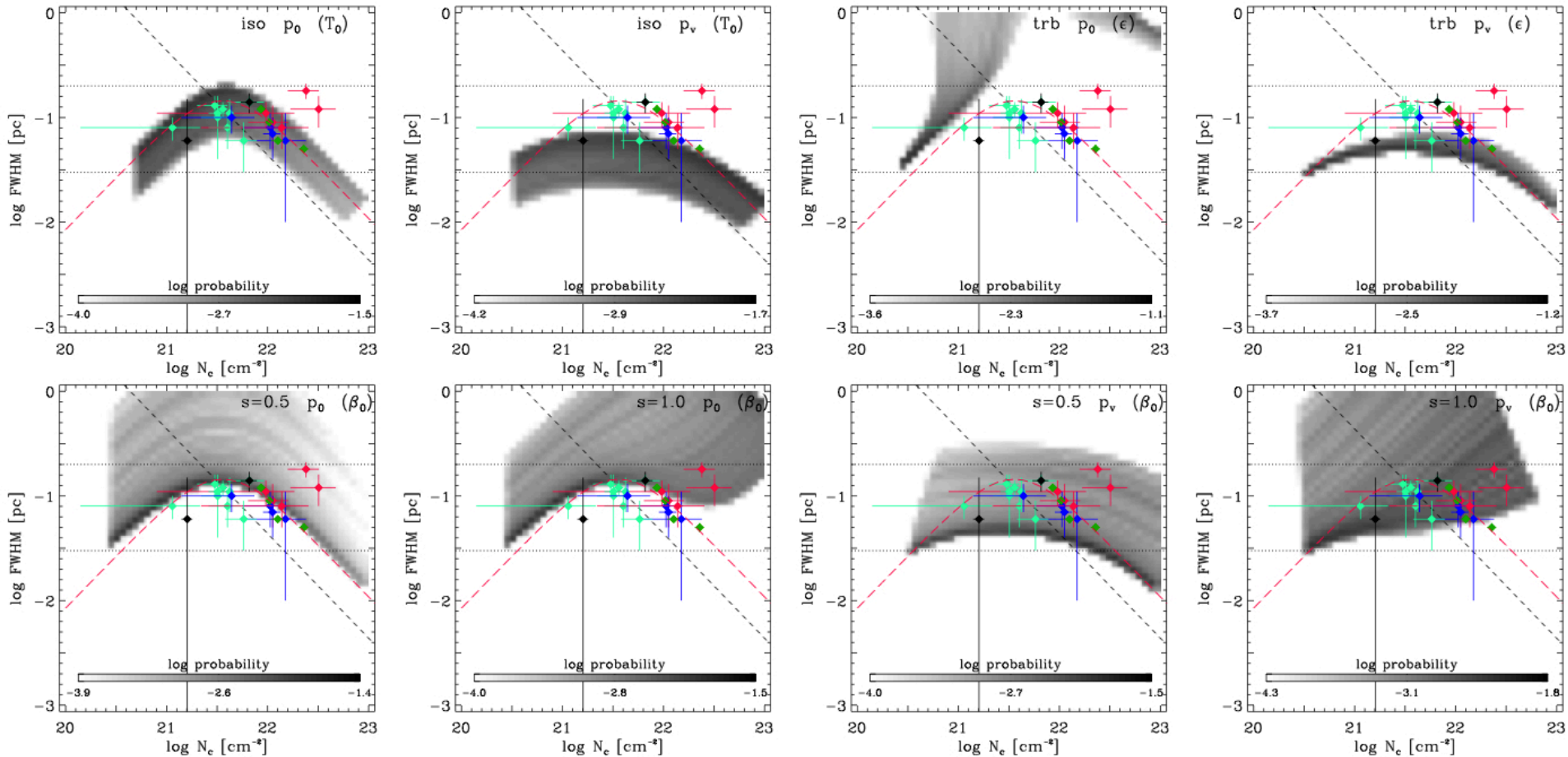


FH13a

(2) Accretion-driven Turbulence

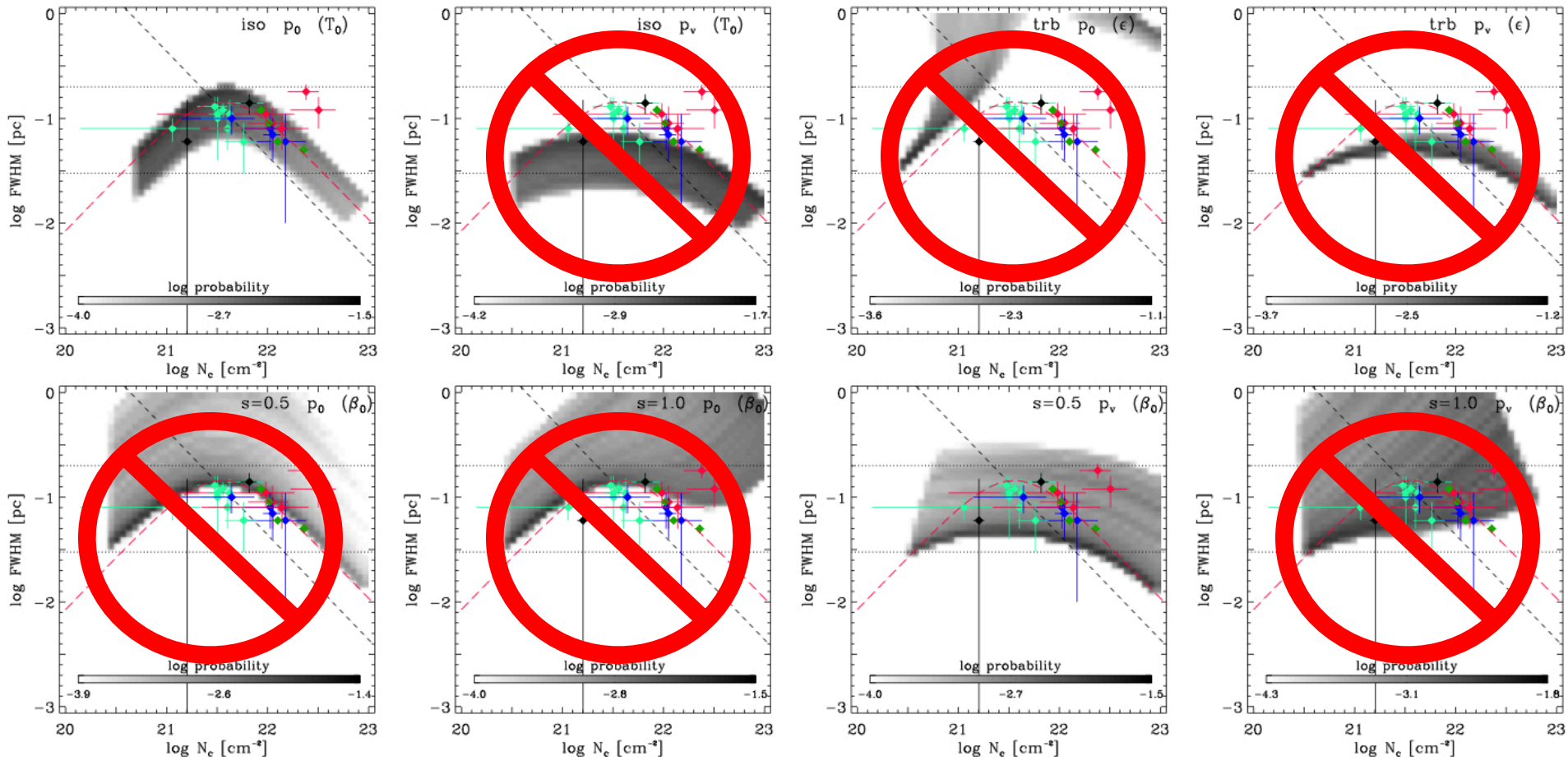
FWHM(N) relations (version 2: $p=2$)

[Fischera & Martin 12a]



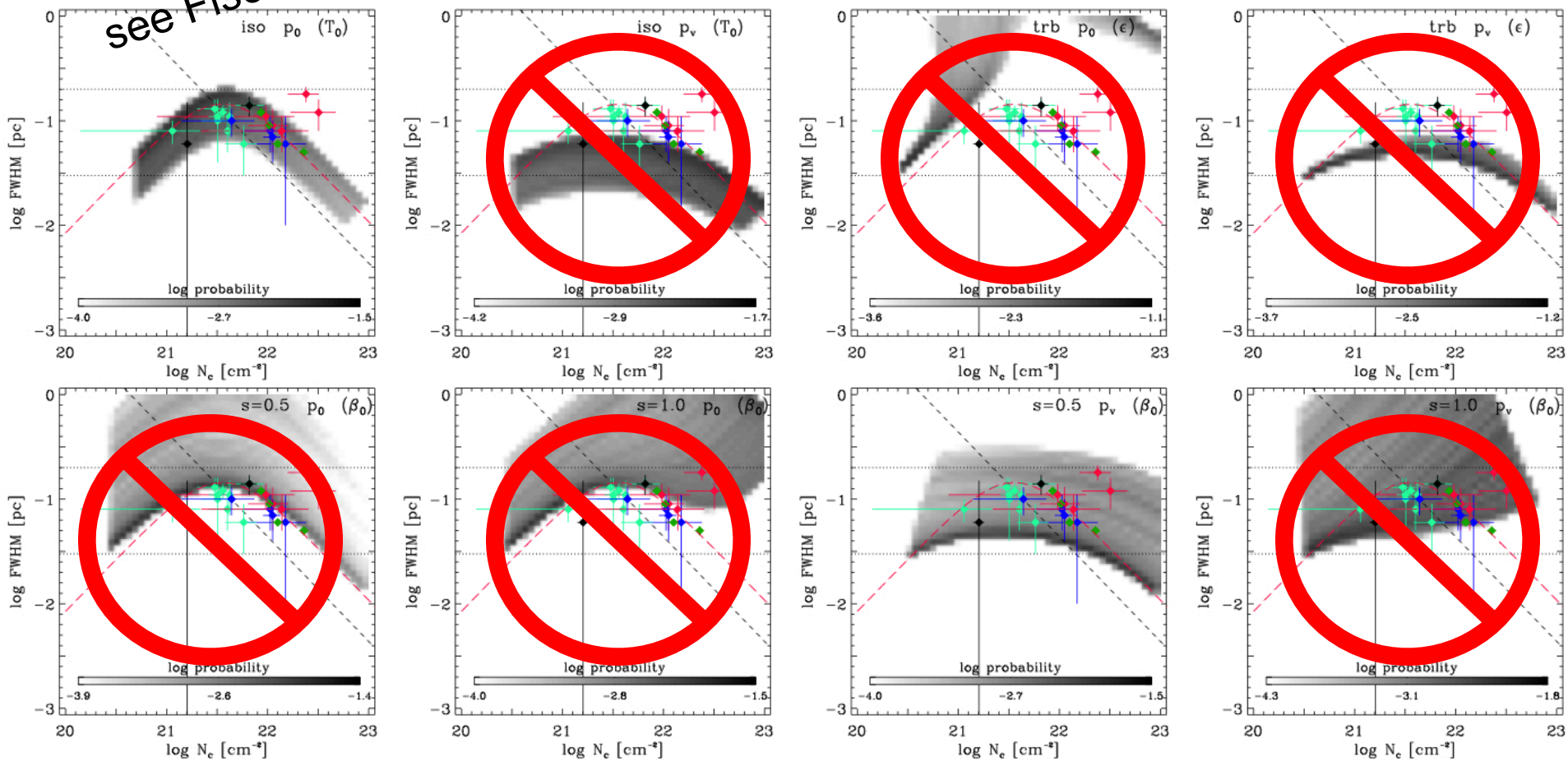
FH13b

(2) Accretion-driven Turbulence



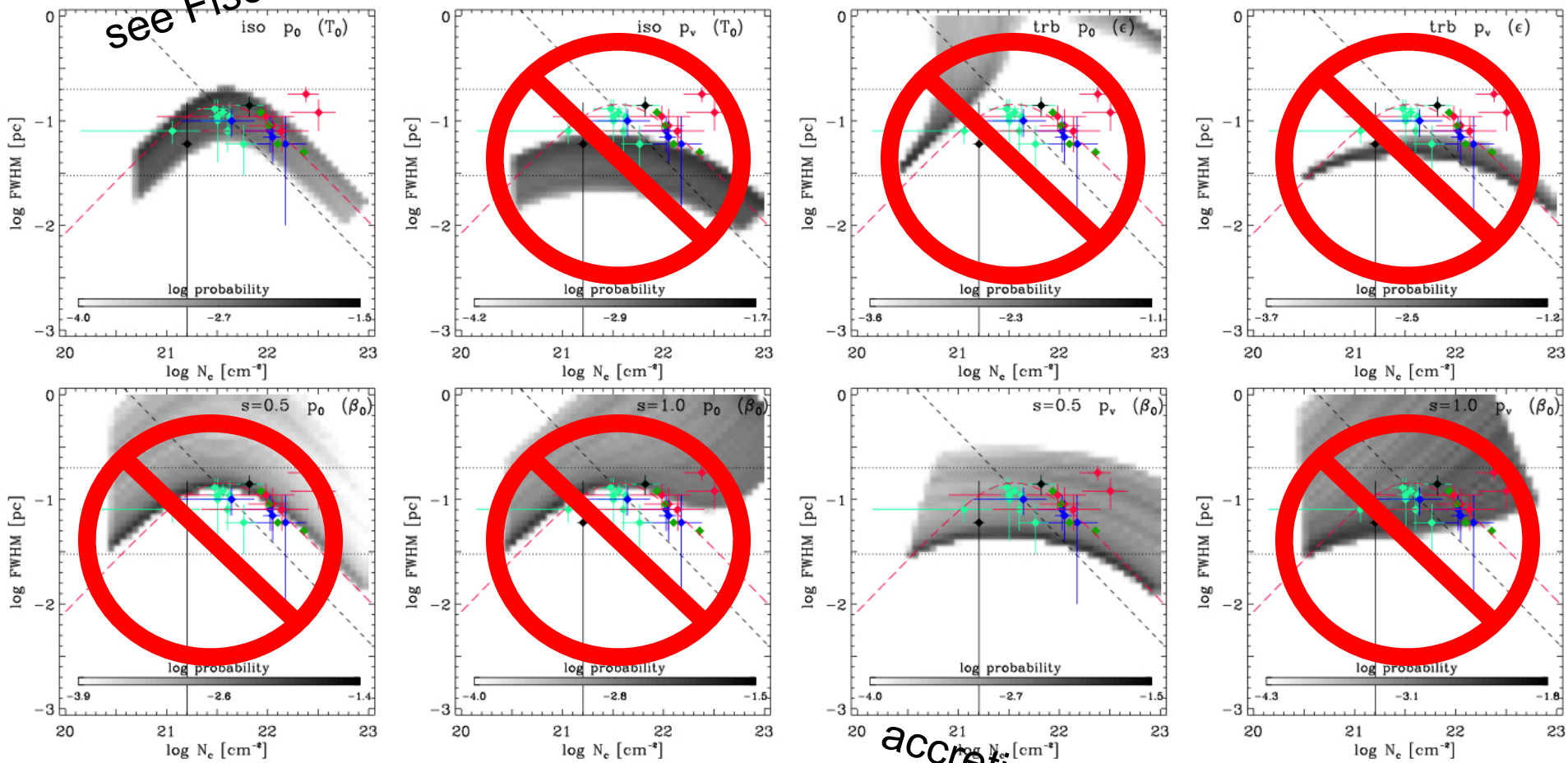
(2) Accretion-driven Turbulence

see Fischera & Martin 12



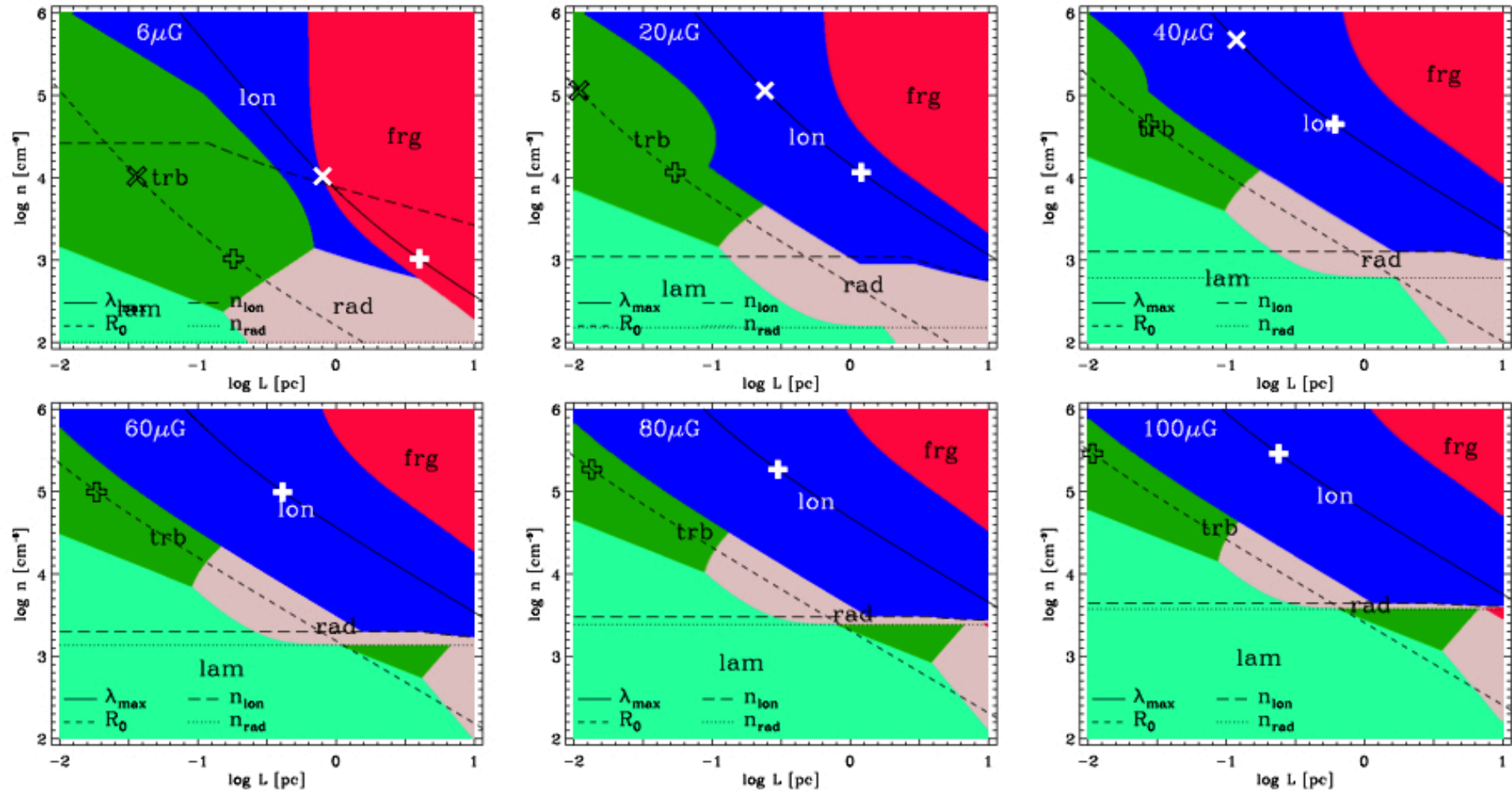
(2) Accretion-driven Turbulence

see Fischera & Martin 12



accretion through layer

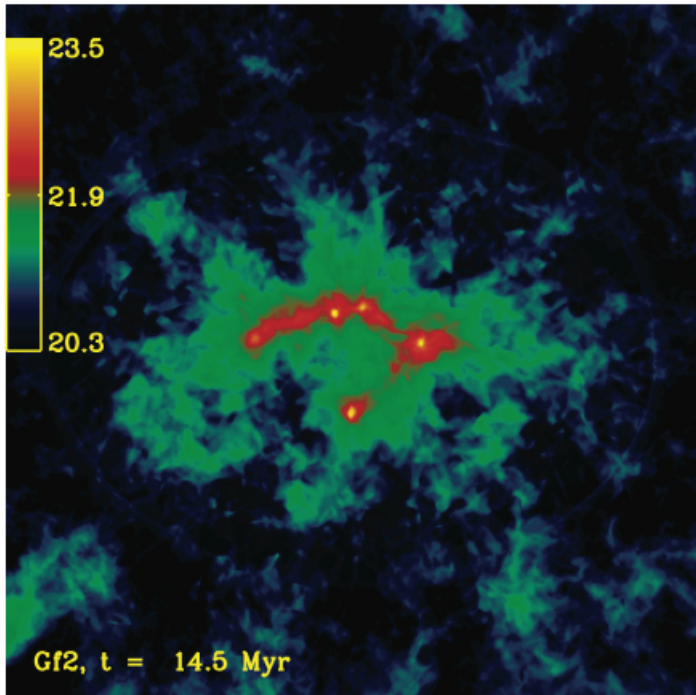
(3) Accretion and Ambipolar Diffusion



FH & Hartmann 14

Filament turns supercritical via accretion, not AD.

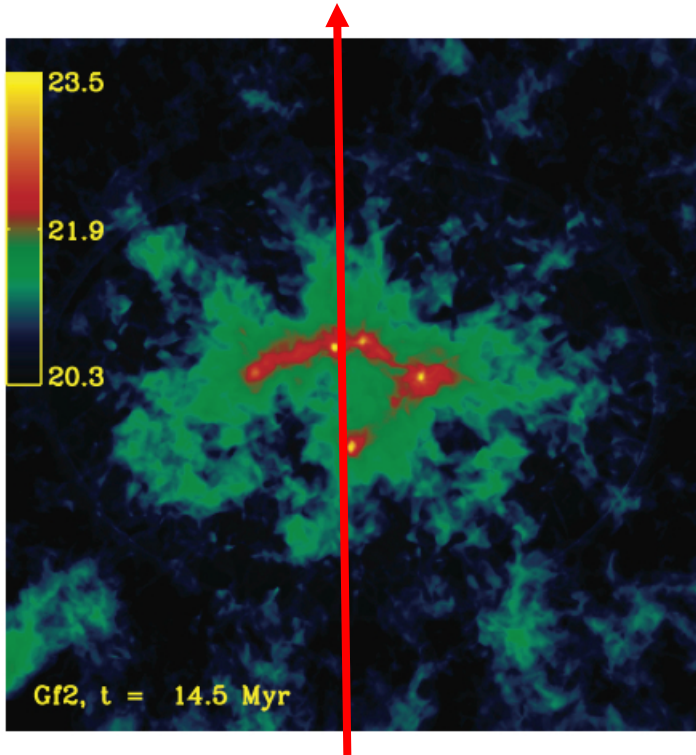
(4) Accretion signatures



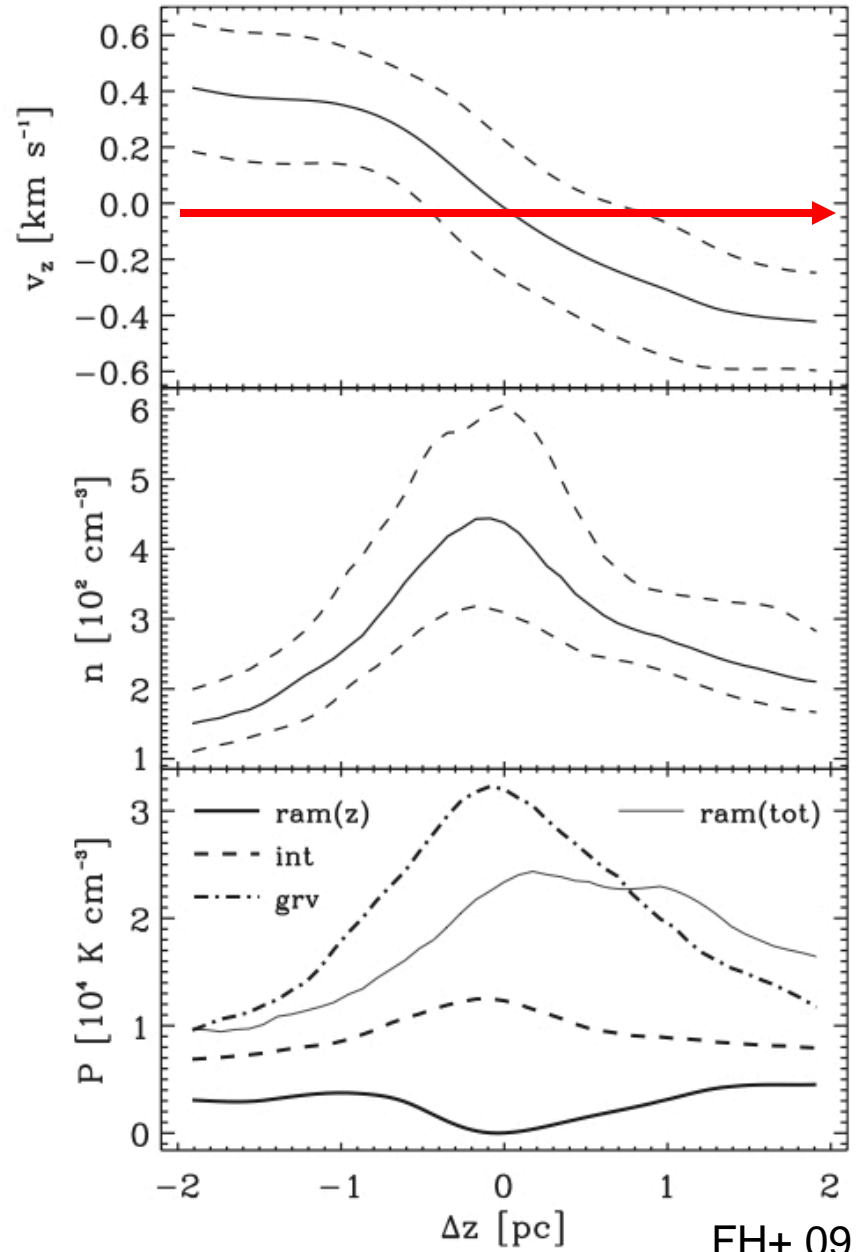
FH+ 08, see Vazquez-Semadeni+ 07,
Banerjee+ 09 etc.

Burkert & Hartmann 04 effect

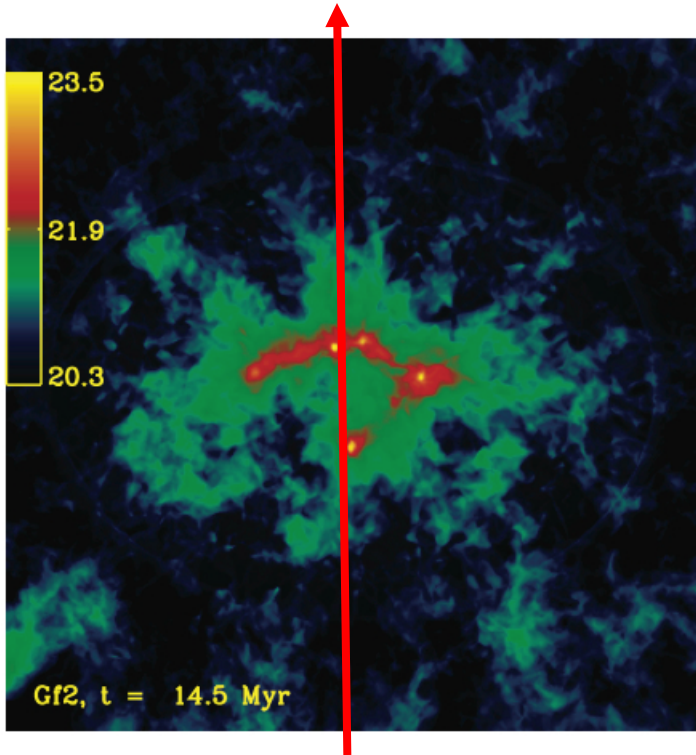
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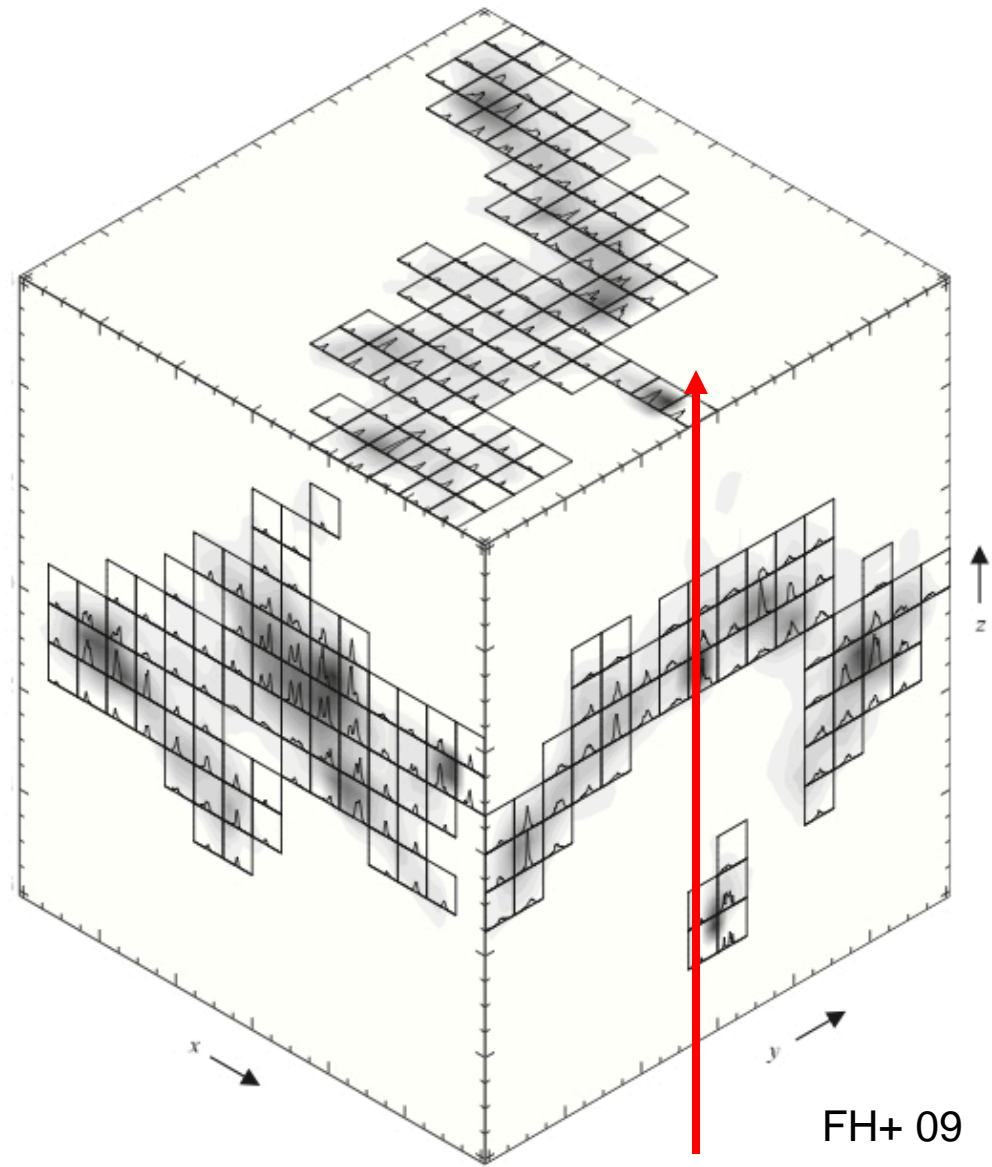
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(4) Accretion signatures

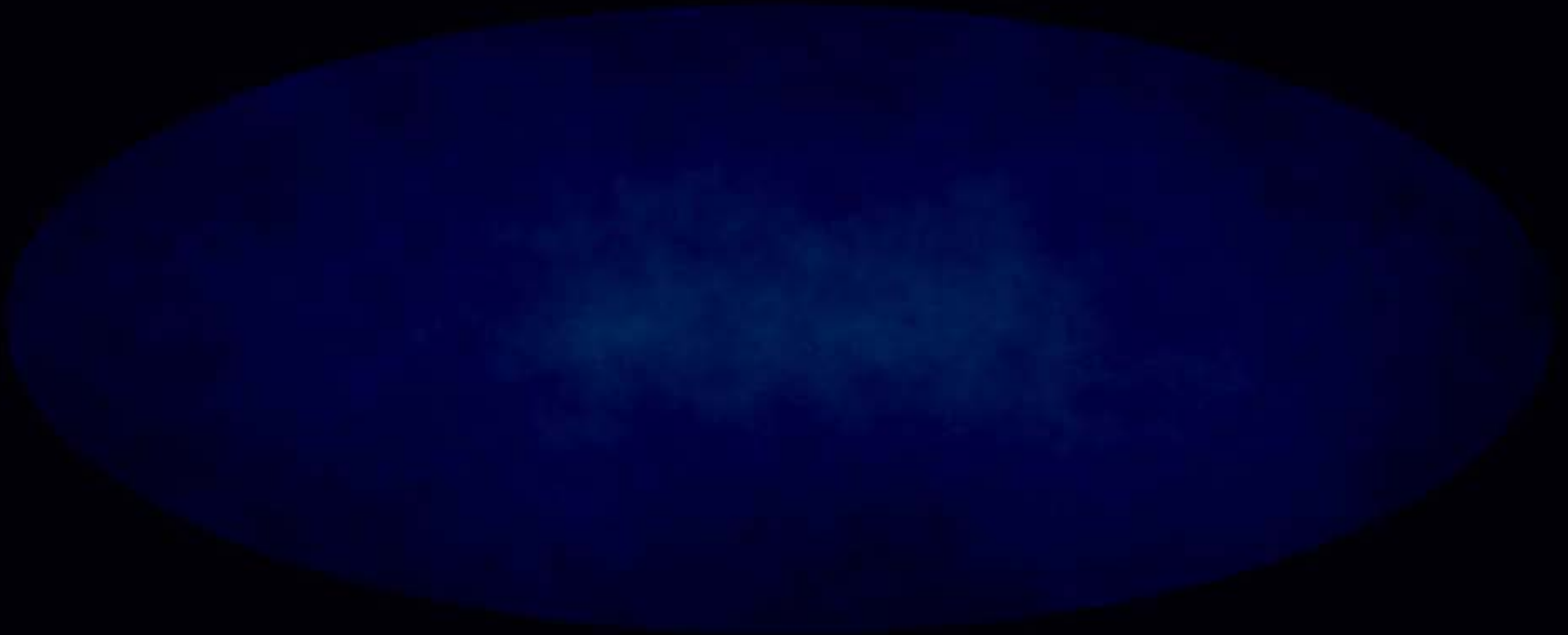


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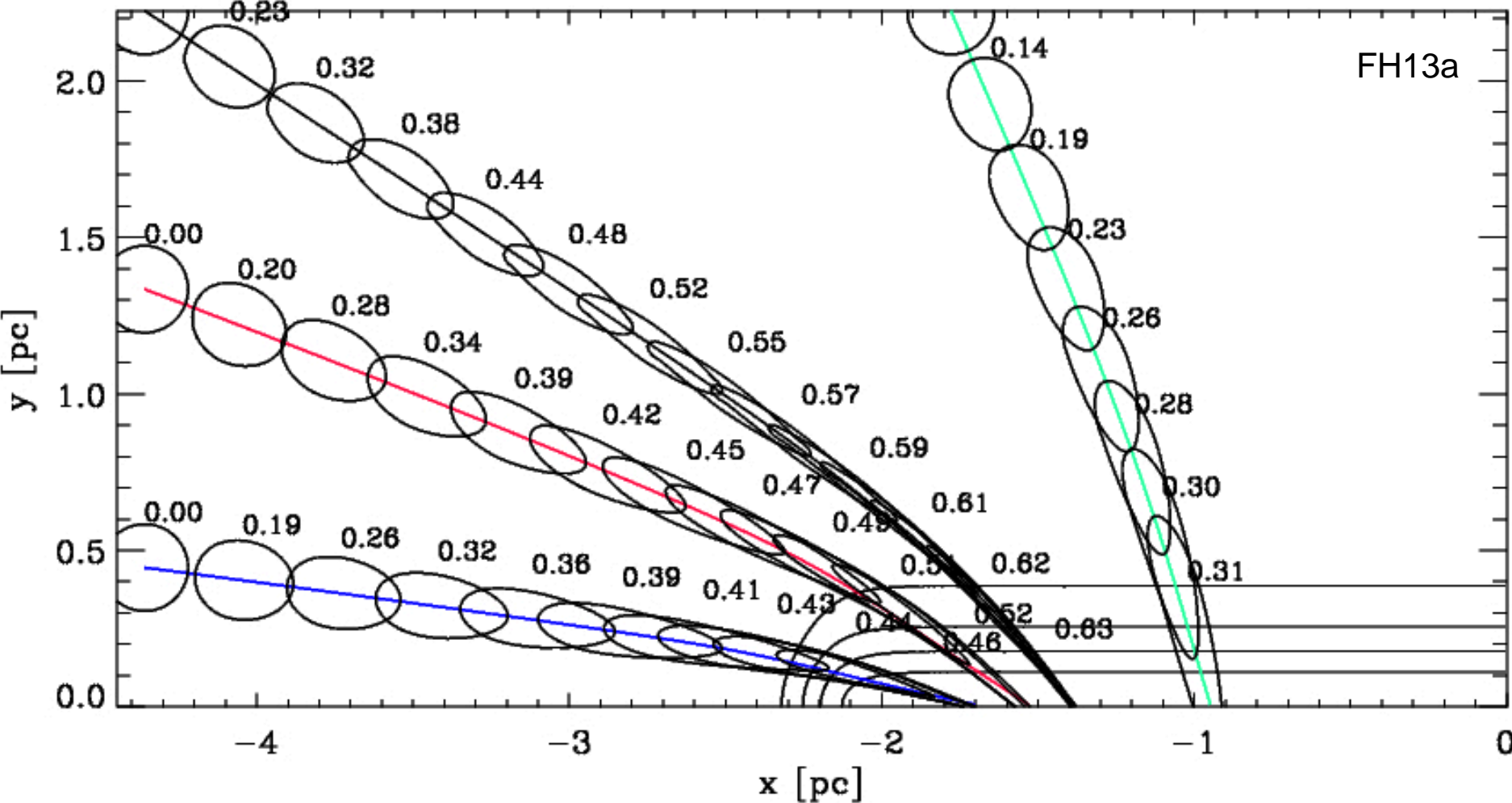


Signature can be hidden in “turbulent” medium.
A filament for us may not be a filament for E.T.

(4) Accretion signatures

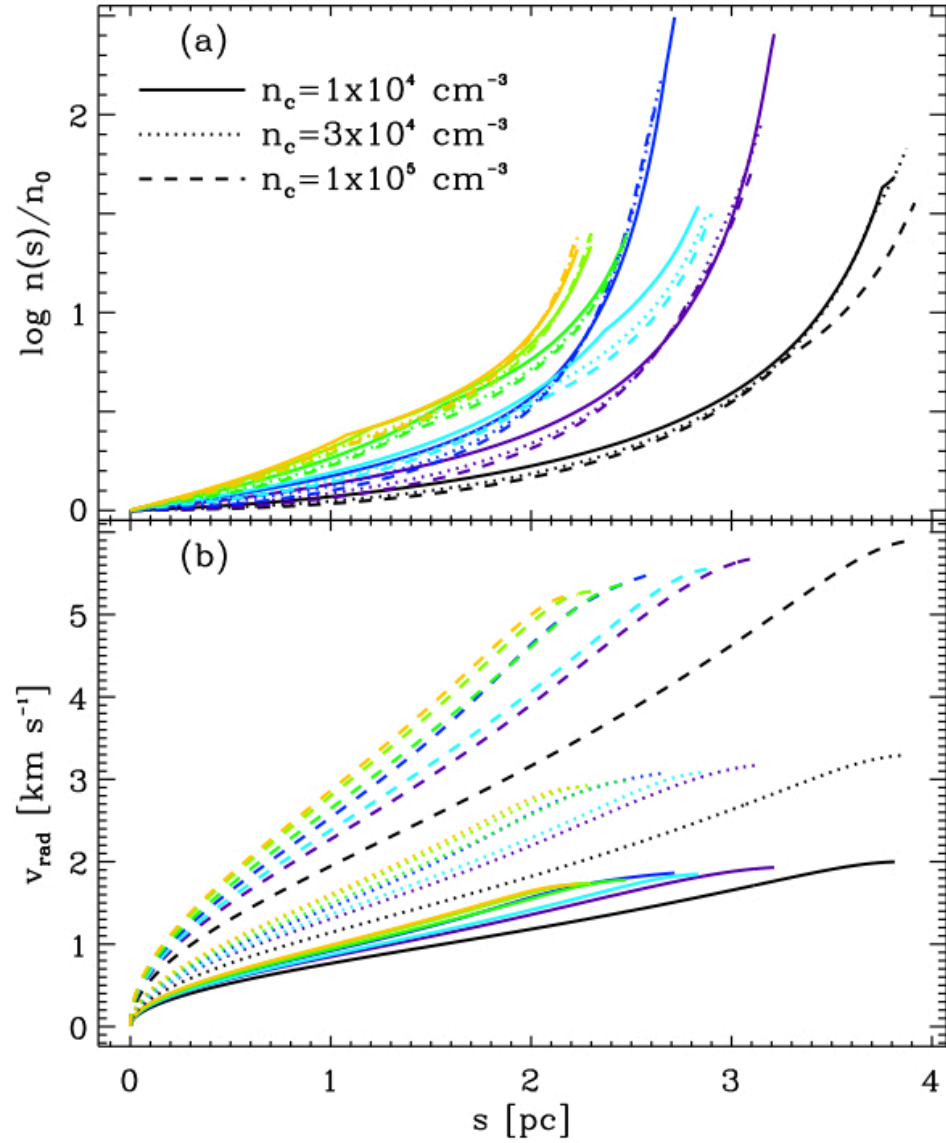
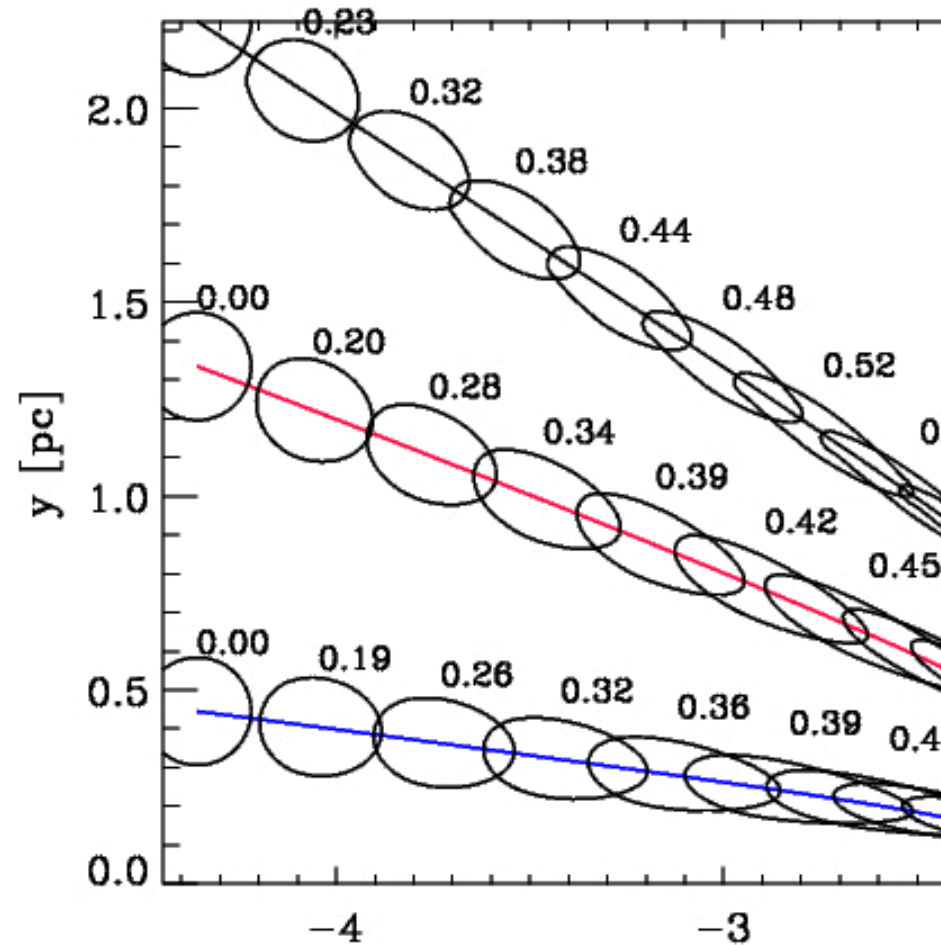


(4) Accretion signatures



Test particle trajectory at filament end

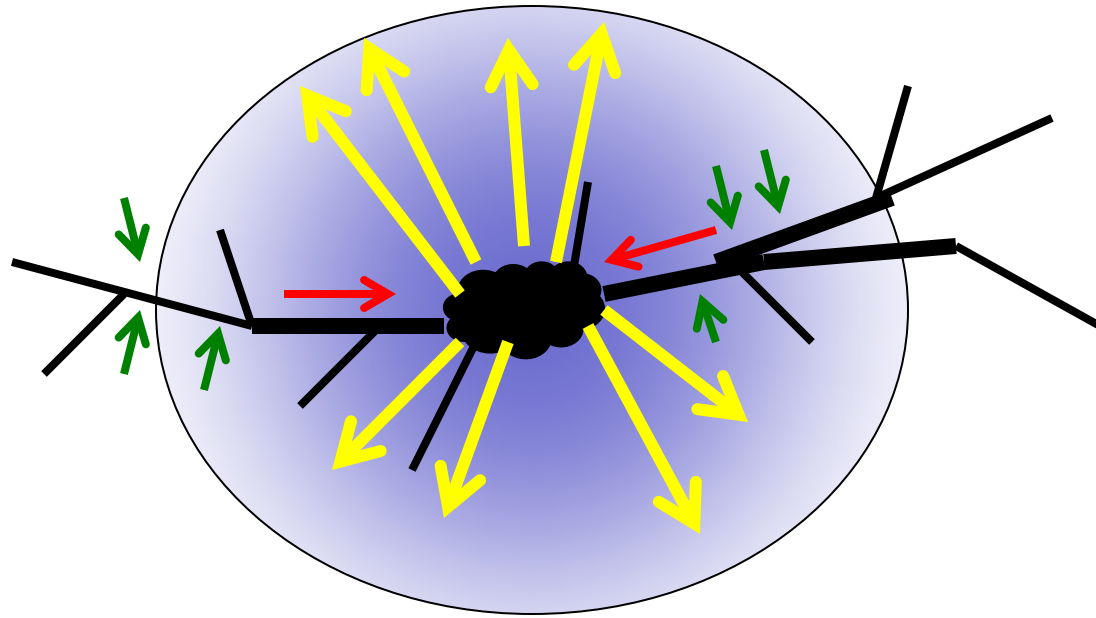
(4) Accretion signatures



Test particle trajectory at filament end



[Spekulatius]



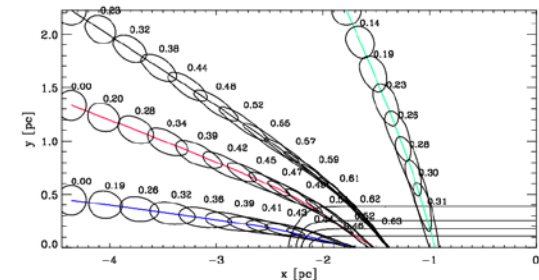
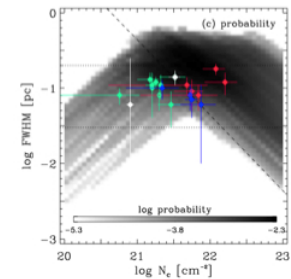
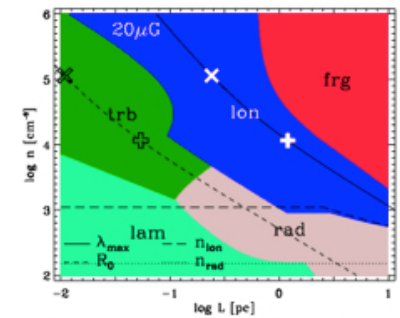
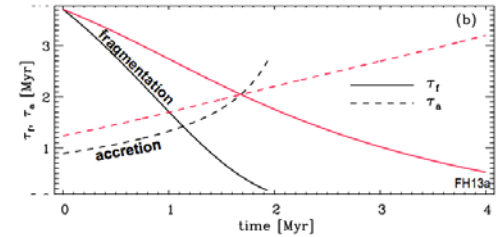
If accretion controls filament evolution, then

(1) longitudinal fragmentation loses.

(2) ambipolar diffusion loses.

(3) “turbulence” is a natural consequence.

(4) its signature could still be hidden in v -space, but revealed in structures.



Thank you!