

NUMERICAL SIMULATIONS
OF
FILAMENT FORMATION
AND FRAGMENTATION

CHE-YU CHEN

WITH EVE OSTRIKER

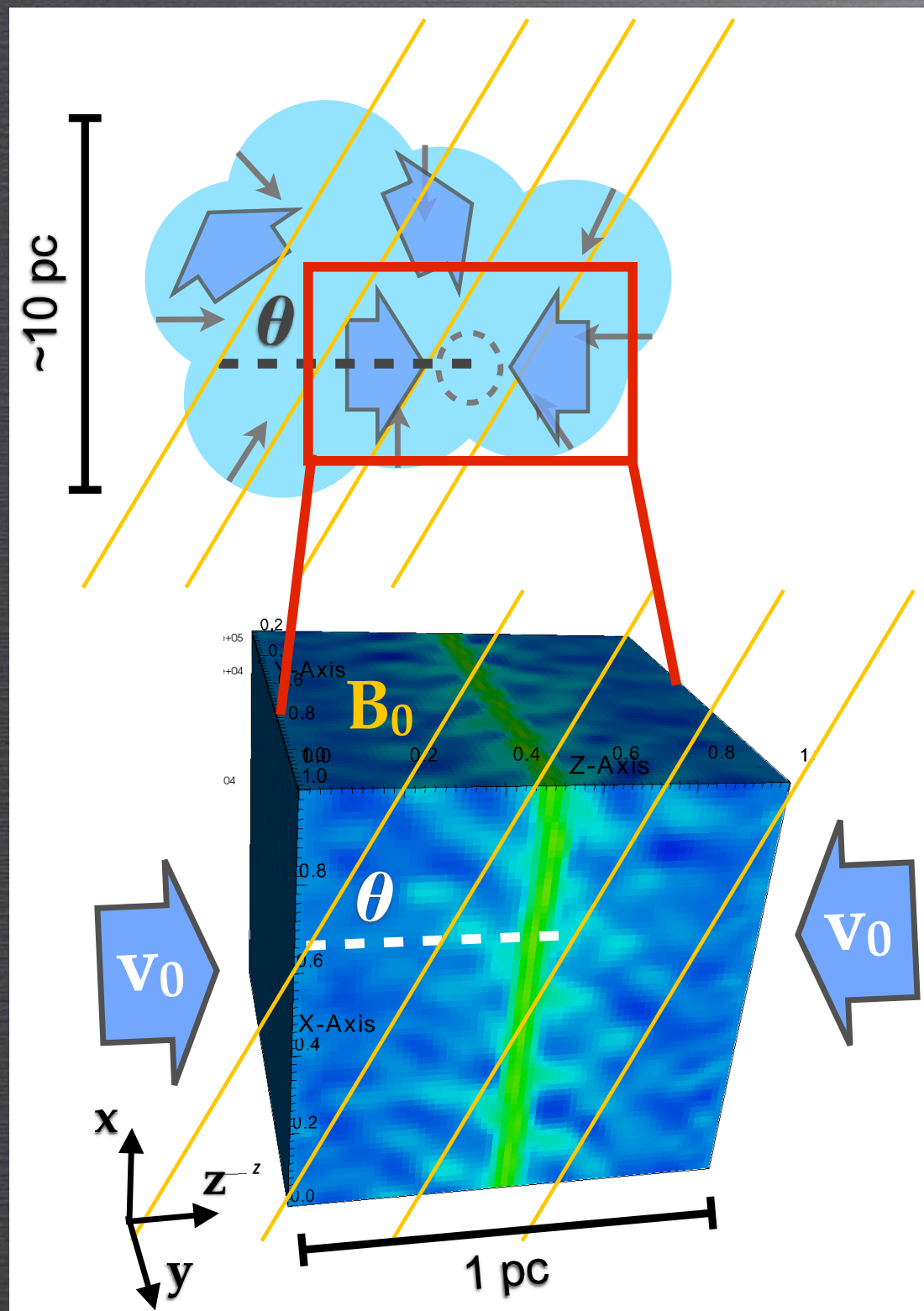
LEE MUNDY

SHAYE STORM

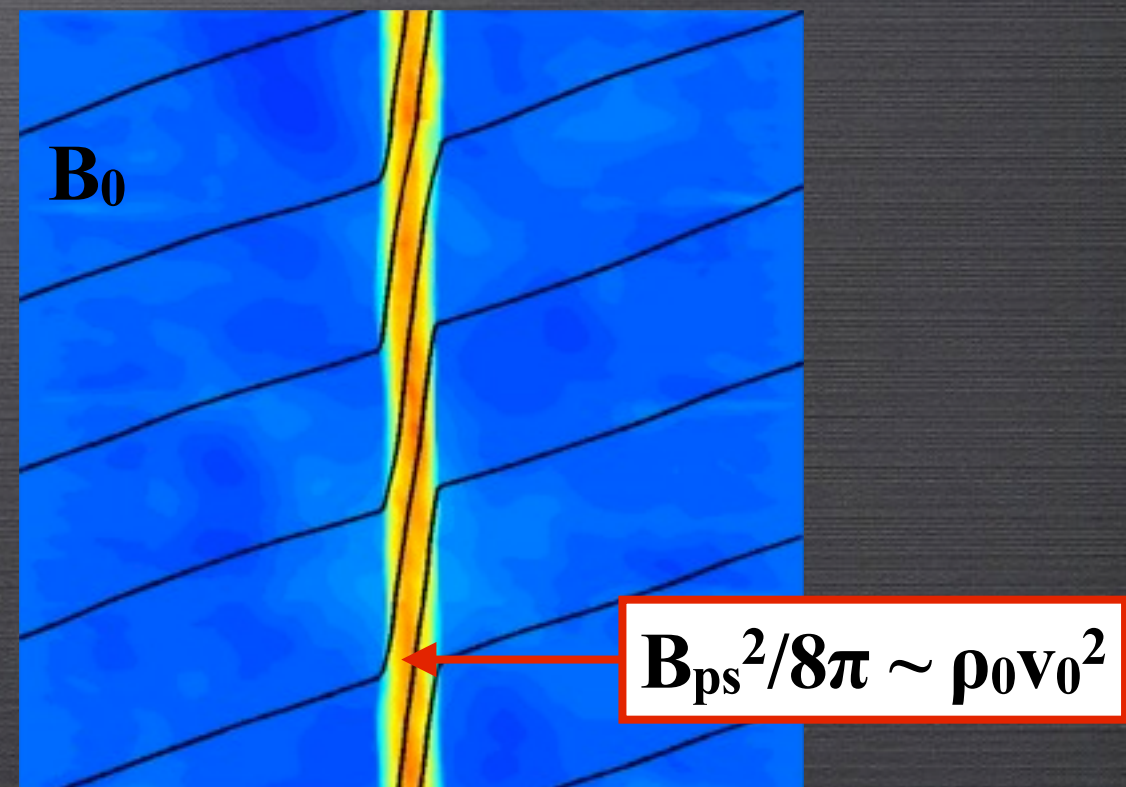
AND THE CLASSY COLLABORATION

FILAMENTS 2014, NRAO

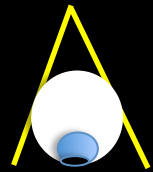
SIMULATION SETUP



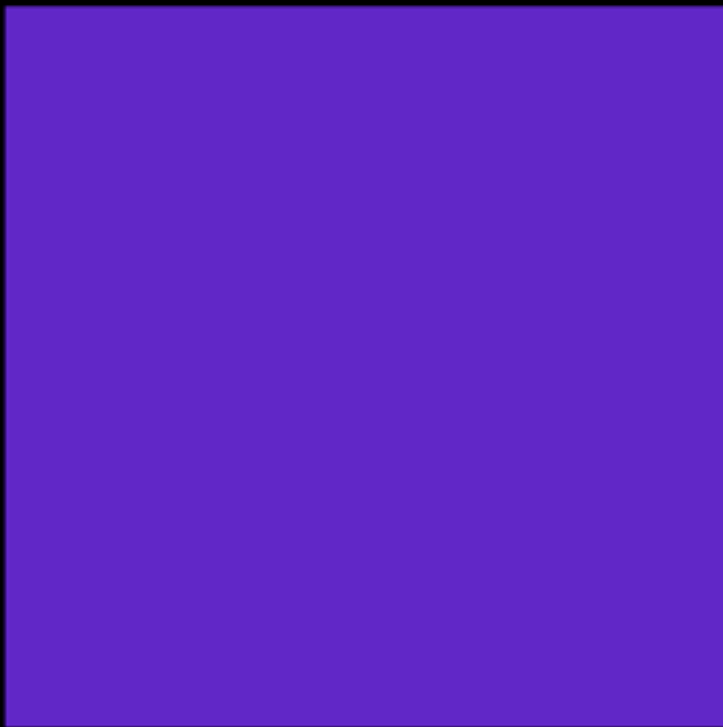
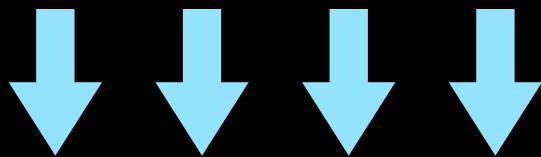
- Include:
 - self-gravity
 - perturbation field
 - ambipolar diffusion
- cloud density:
 $n_0 = 1000 \text{ cm}^{-3}$



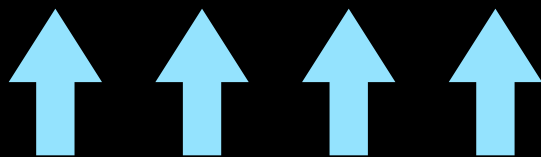
STRUCTURE FORMATION



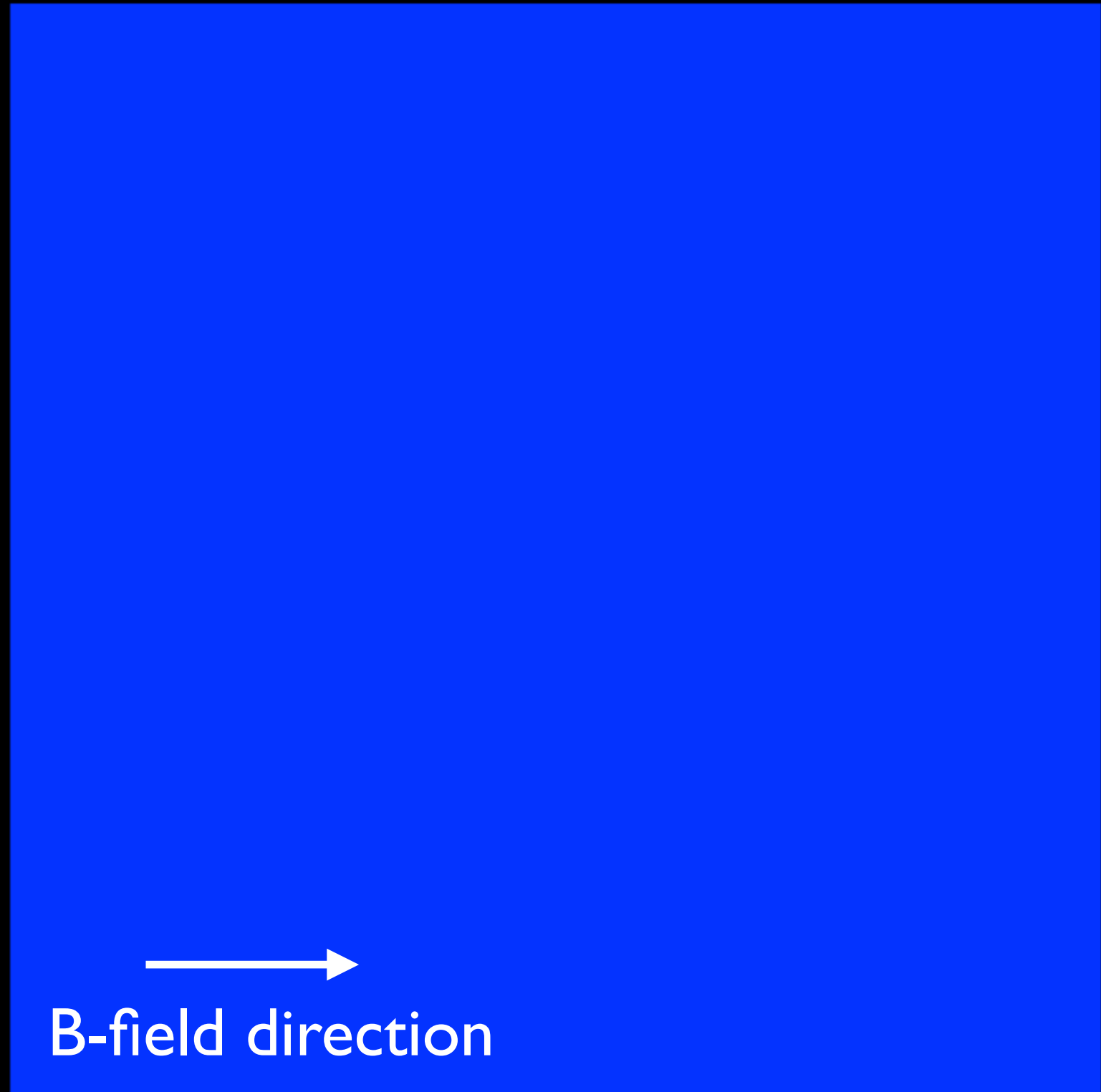
inflow



$t = 0.000 \text{ Myr}$



inflow

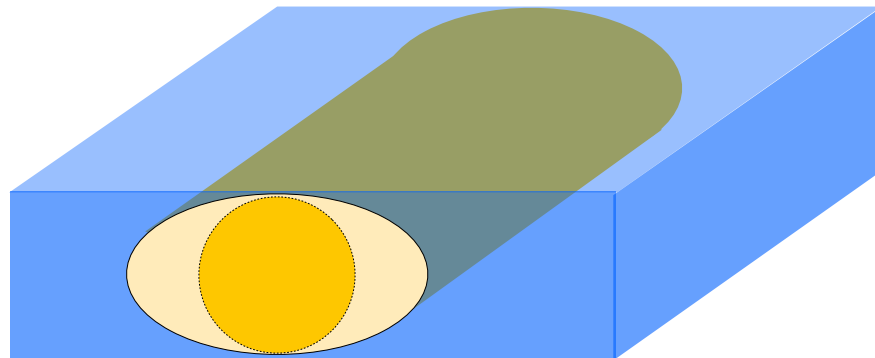


B-field direction

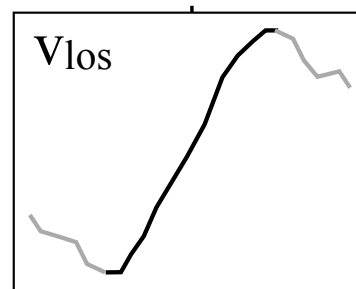
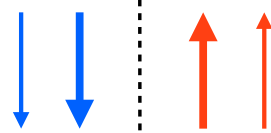
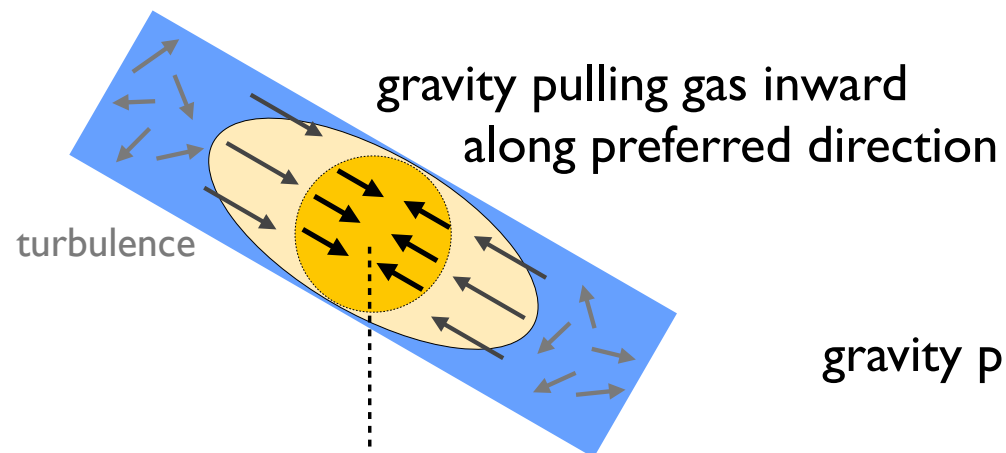
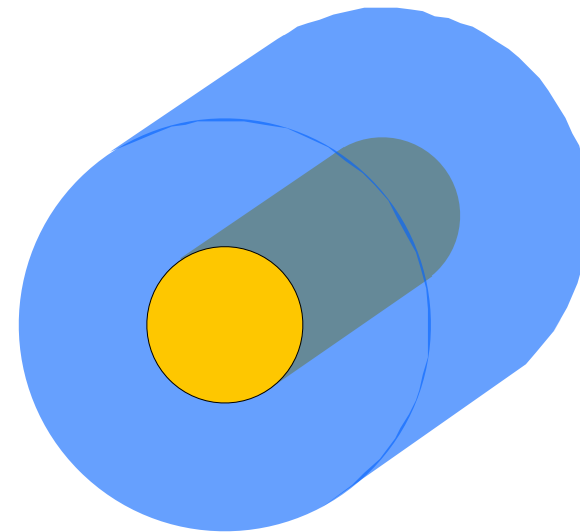
$t = 0.000 \text{ Myr}$

KINEMATICS

filament forming in a slab

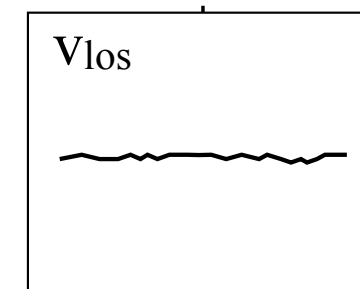
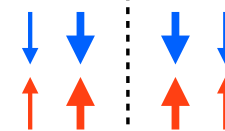
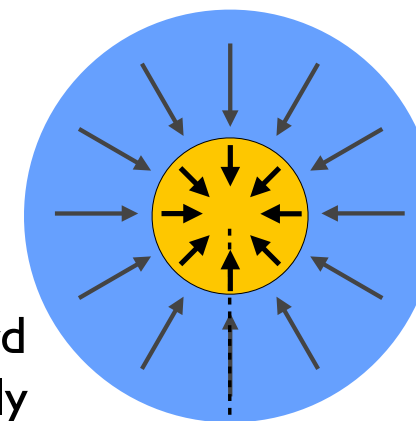


filament forming in a cylinder



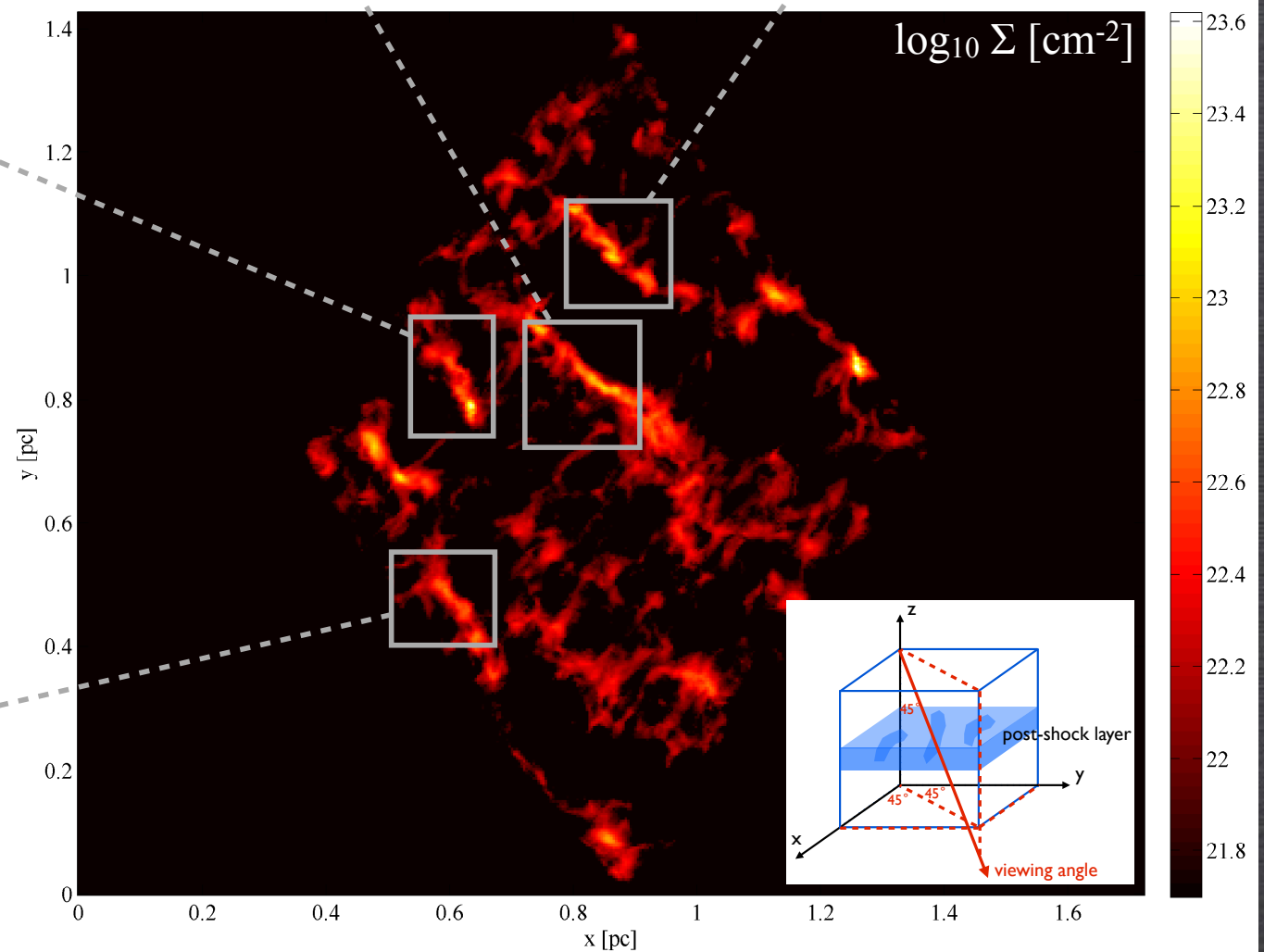
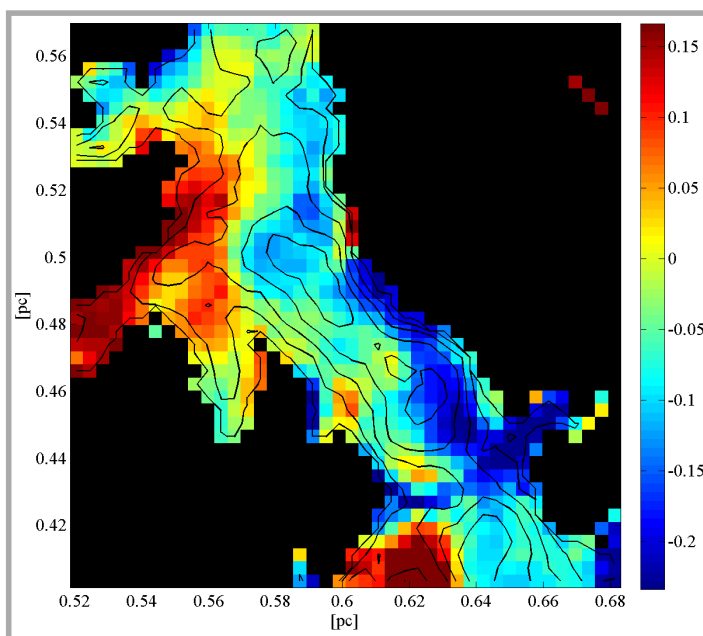
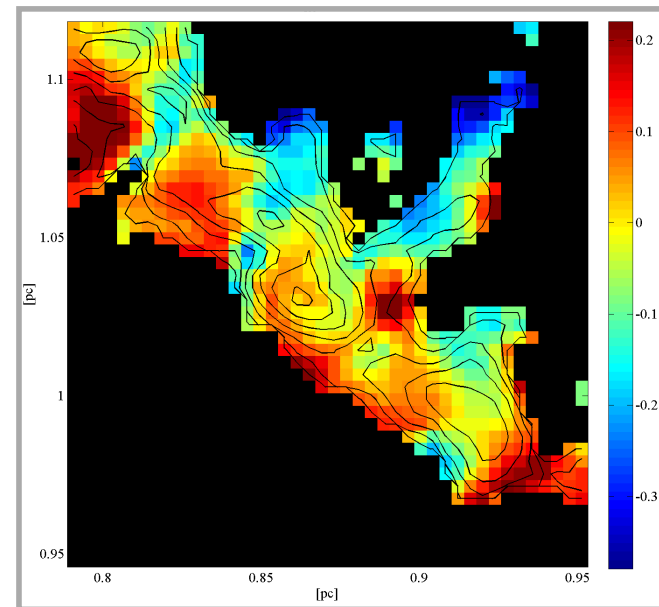
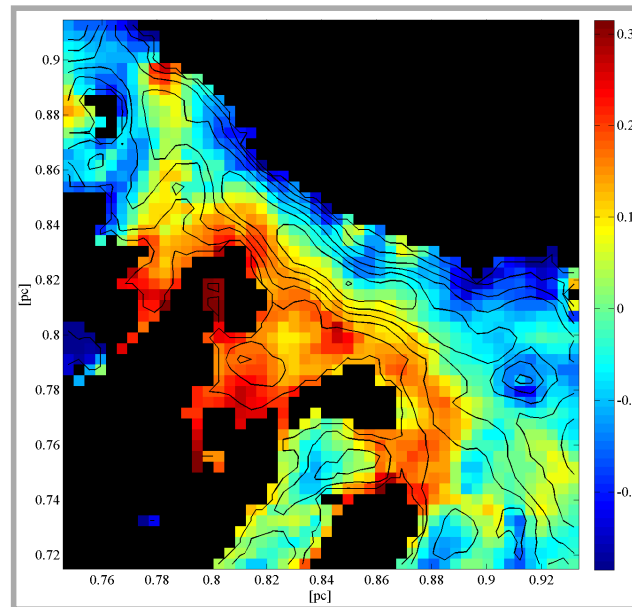
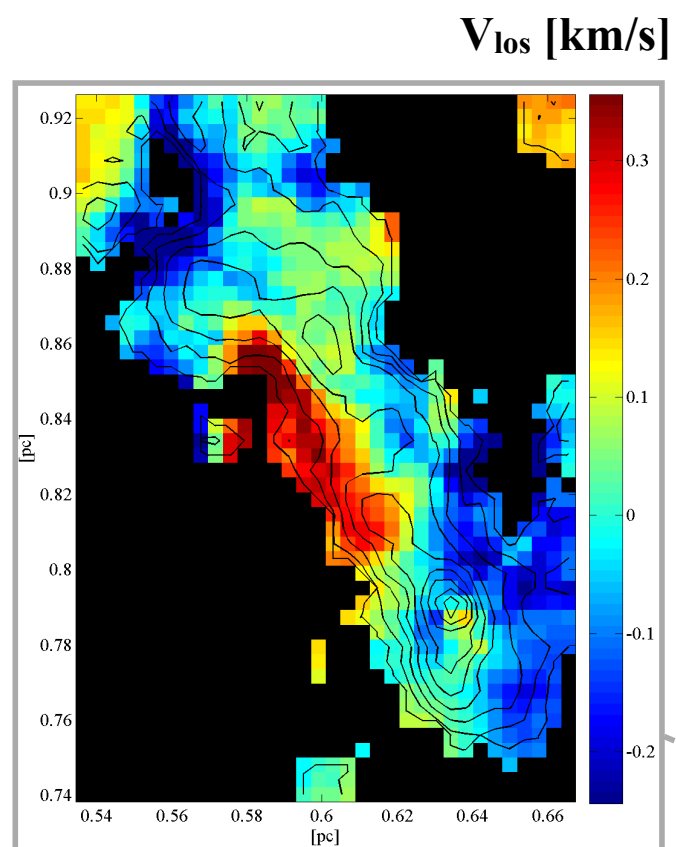
observer

gravity pulling gas inward isotropically



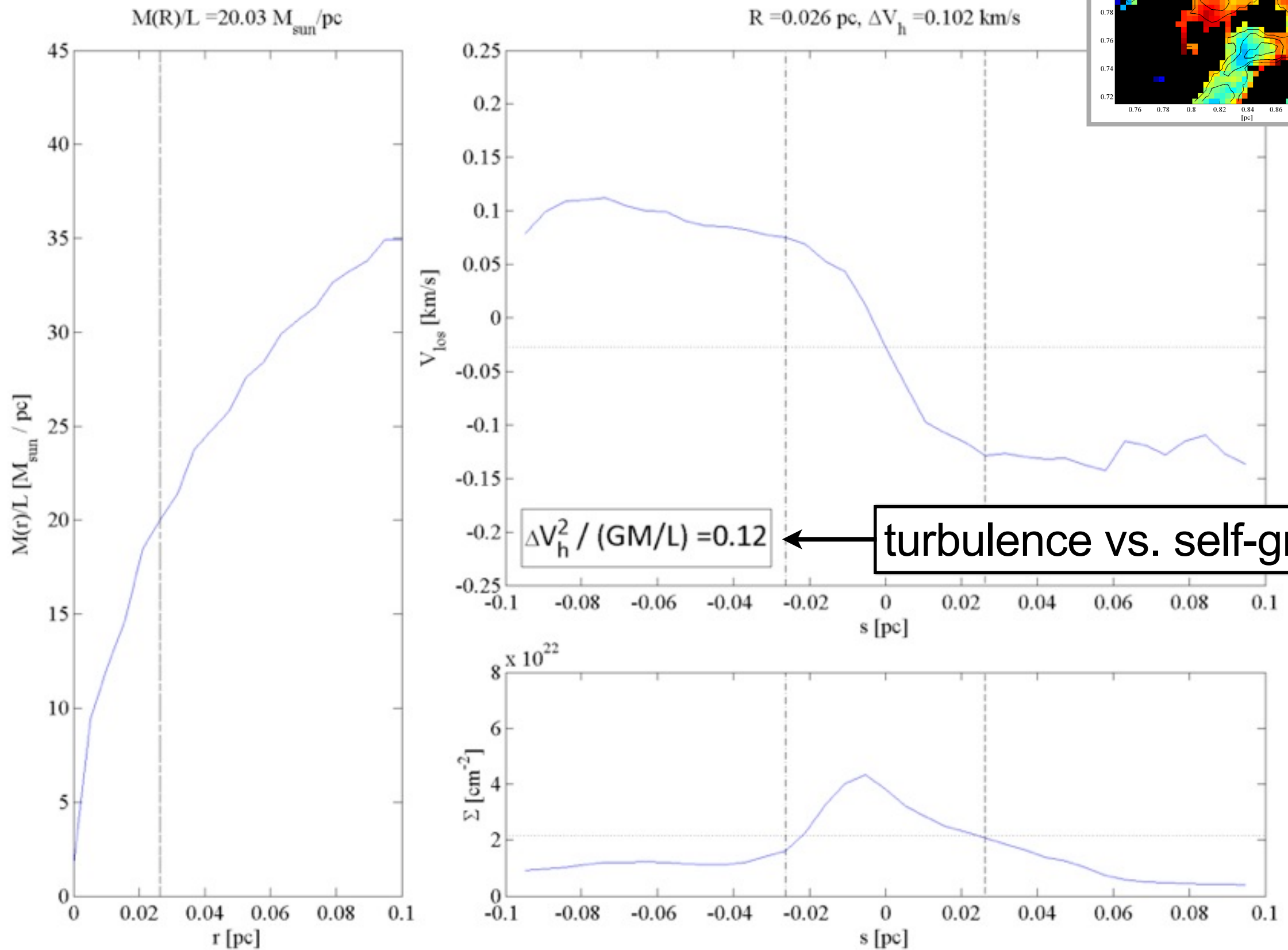
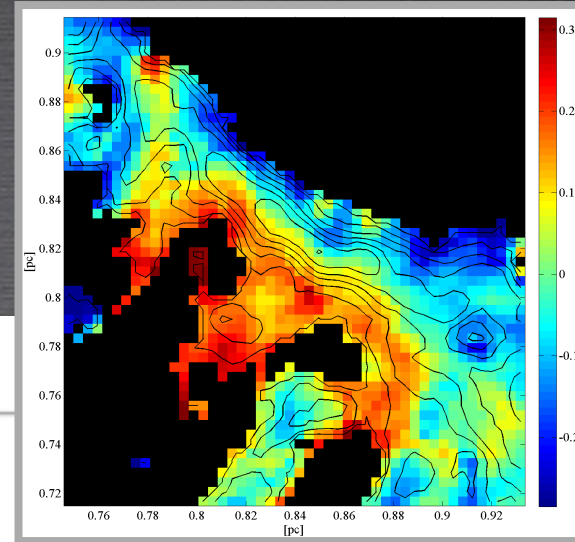
observer

LINE-OF-SIGHT VELOCITY



(Mundy et al. *in prep.*)

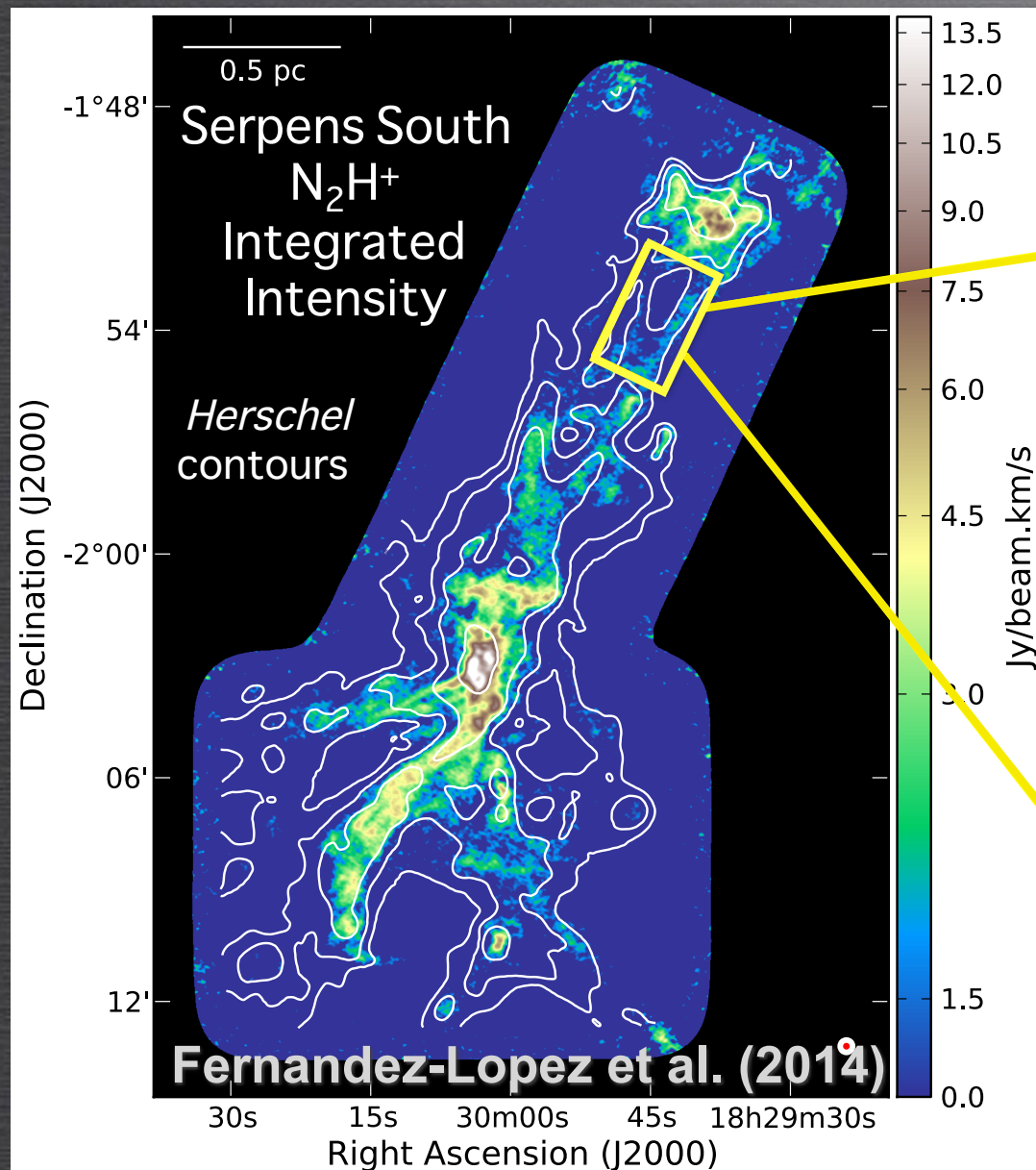
FILAMENT STRUCTURE



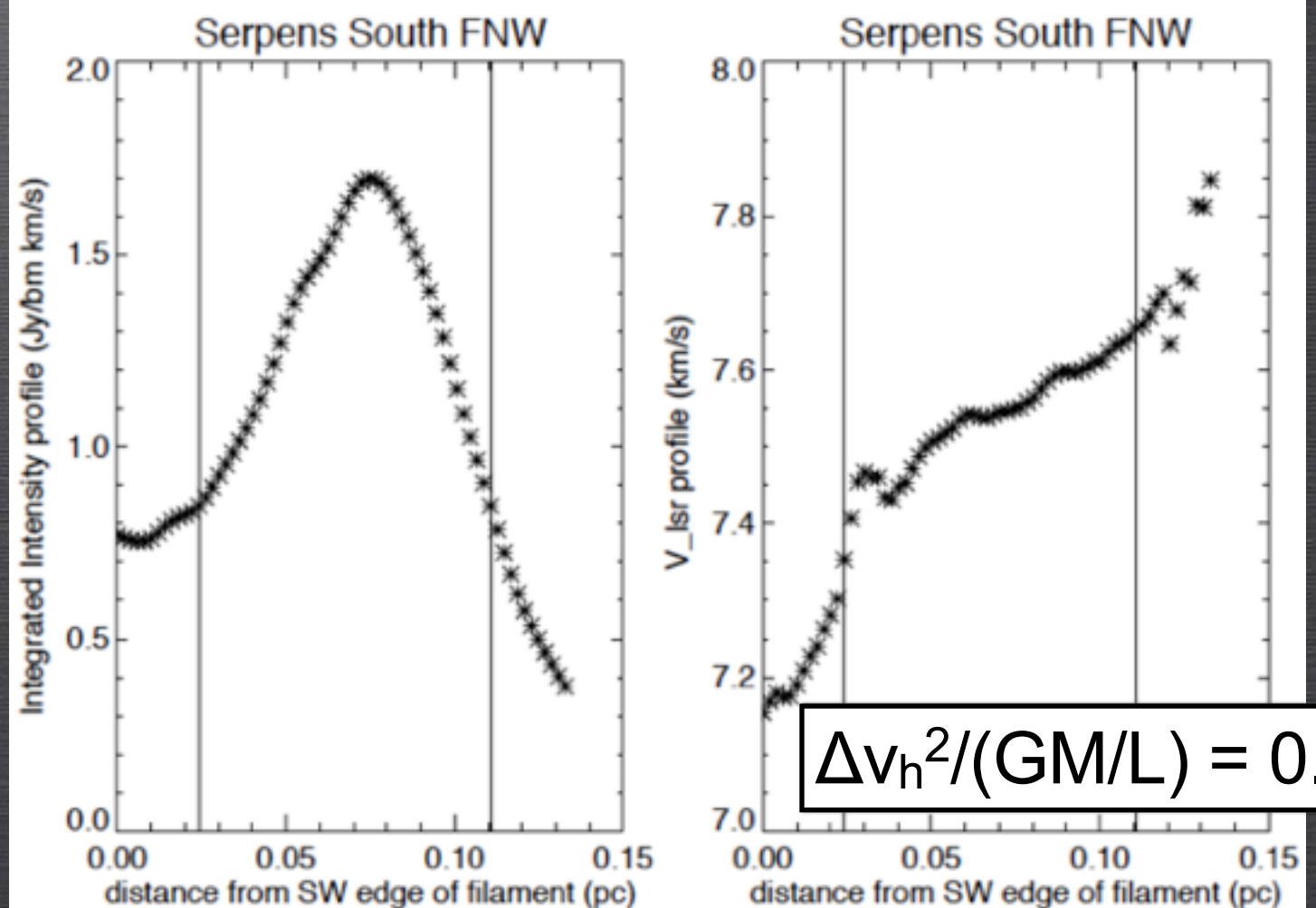
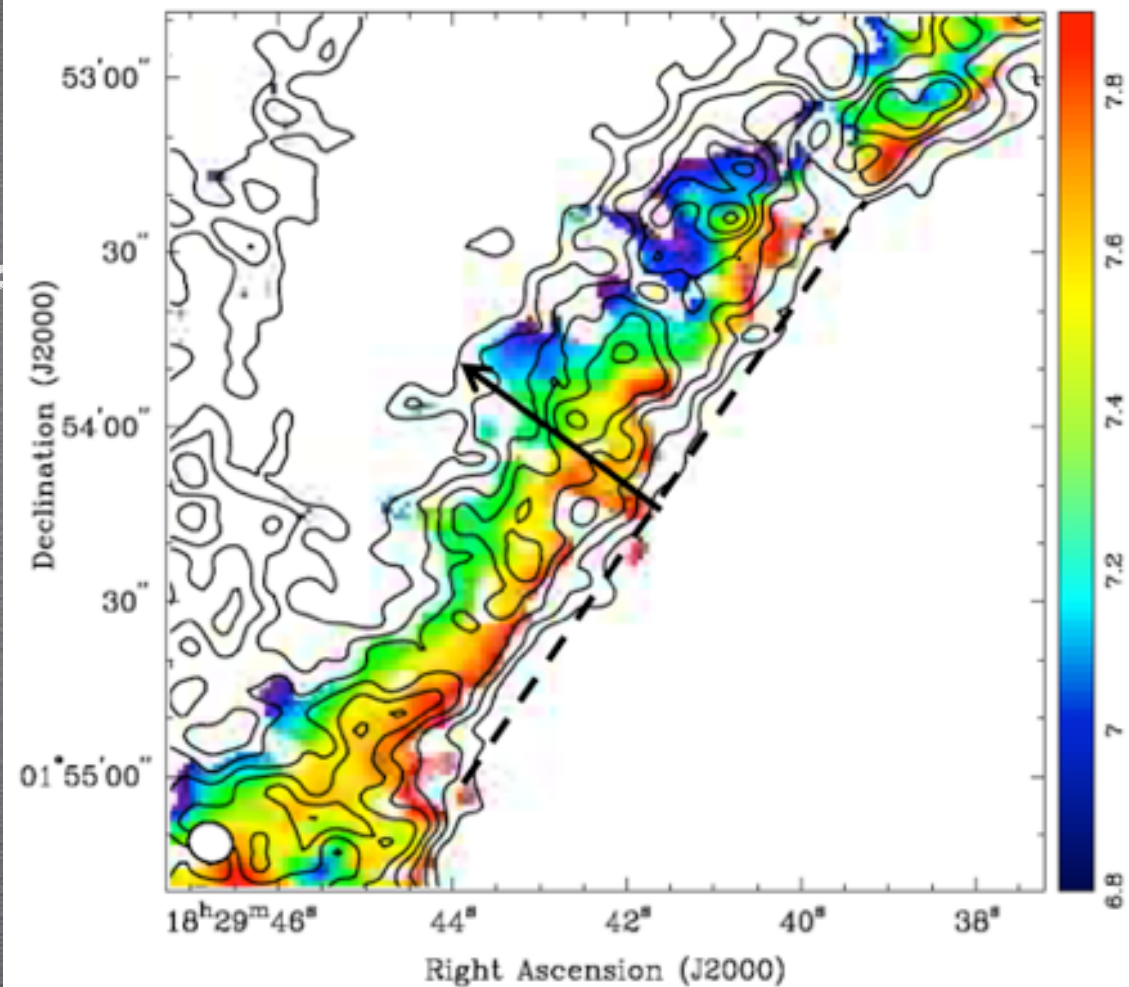
(Mundy et al. *in prep.*)

OBSERVATIONS

- CLASSy results

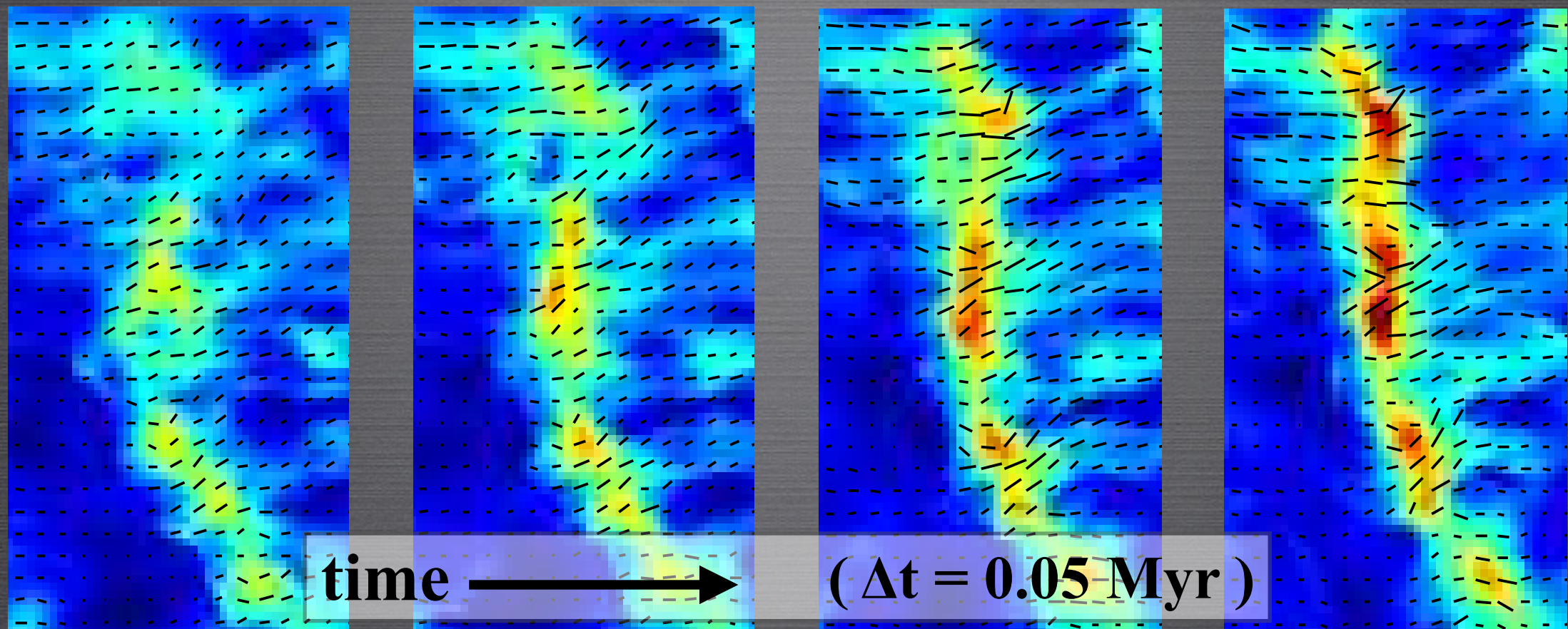


(Mundy et al. in prep.)

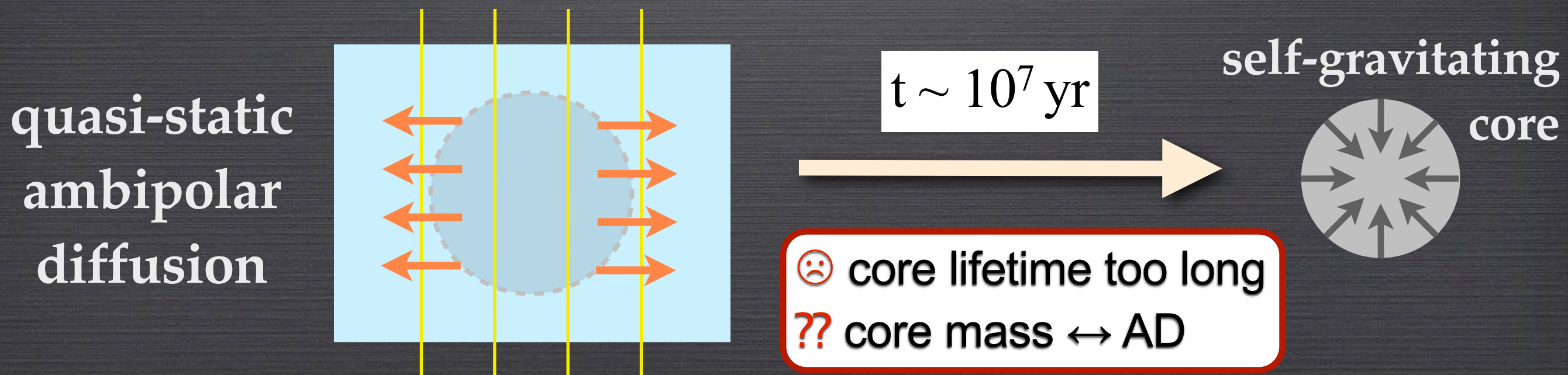


see also S. Storm's poster
K. Lee's poster

FILAMENT $\Rightarrow$ CORE FORMATION

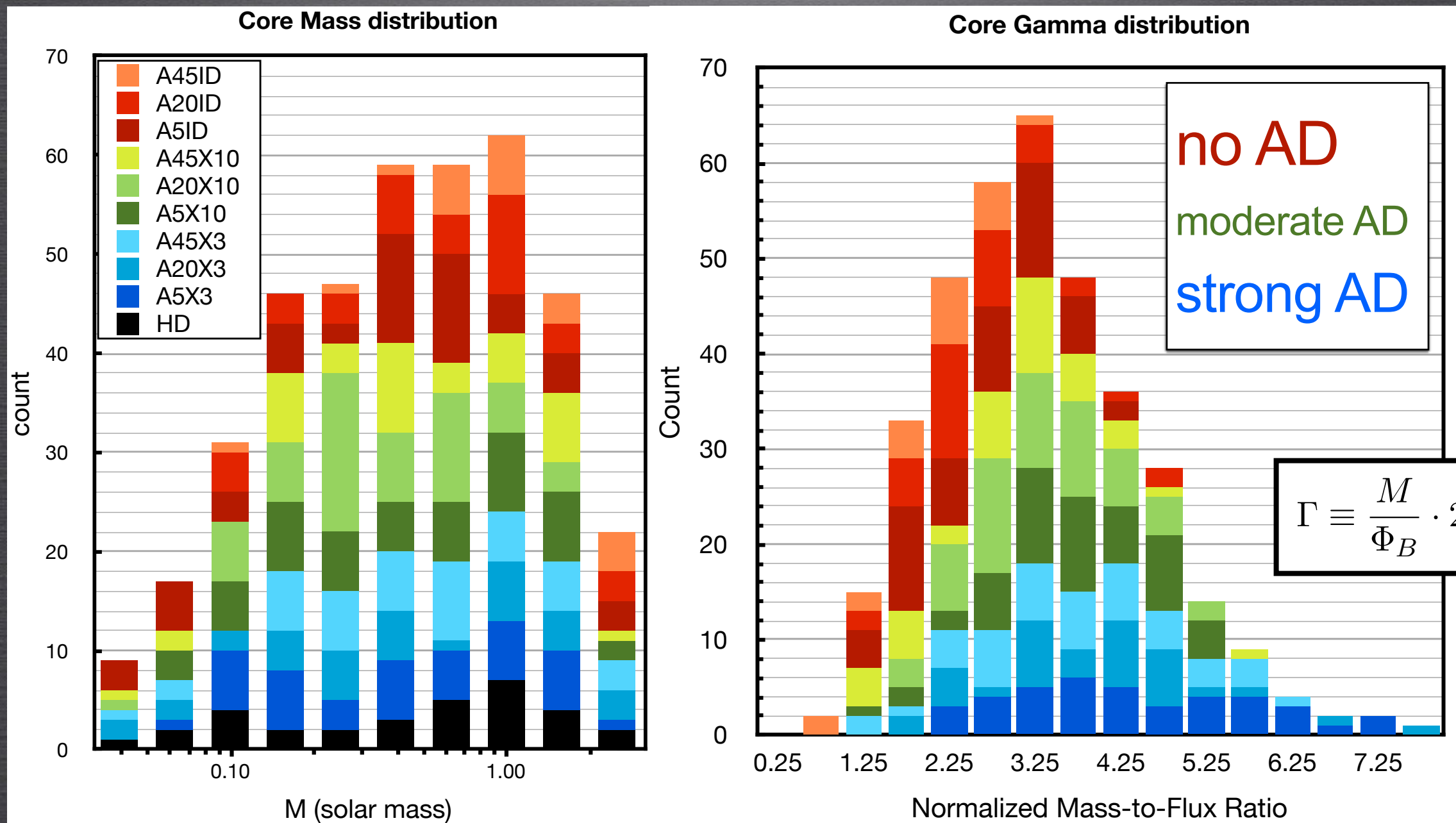


- Classical picture of magnetized cores

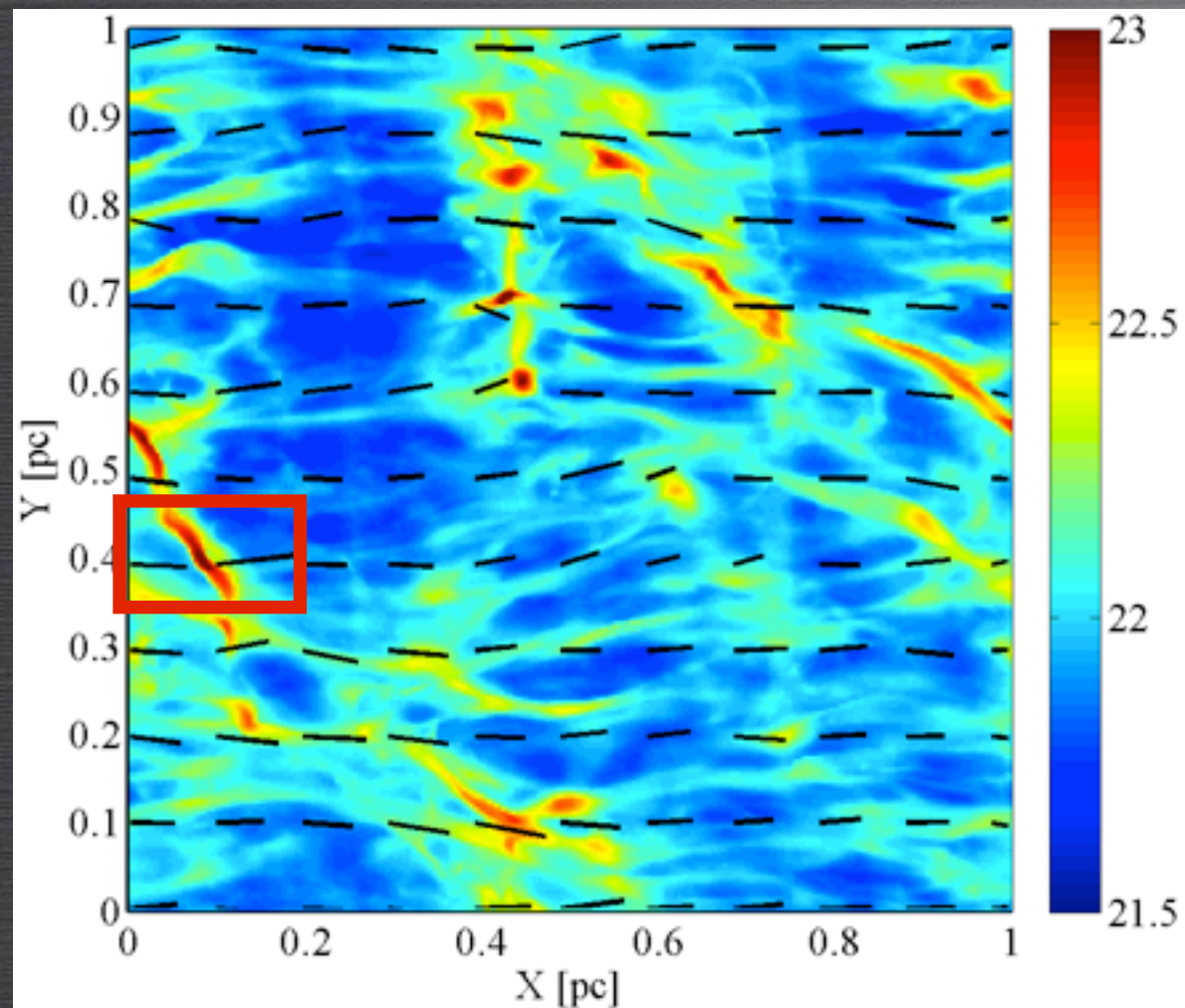


SIMULATION RESULTS

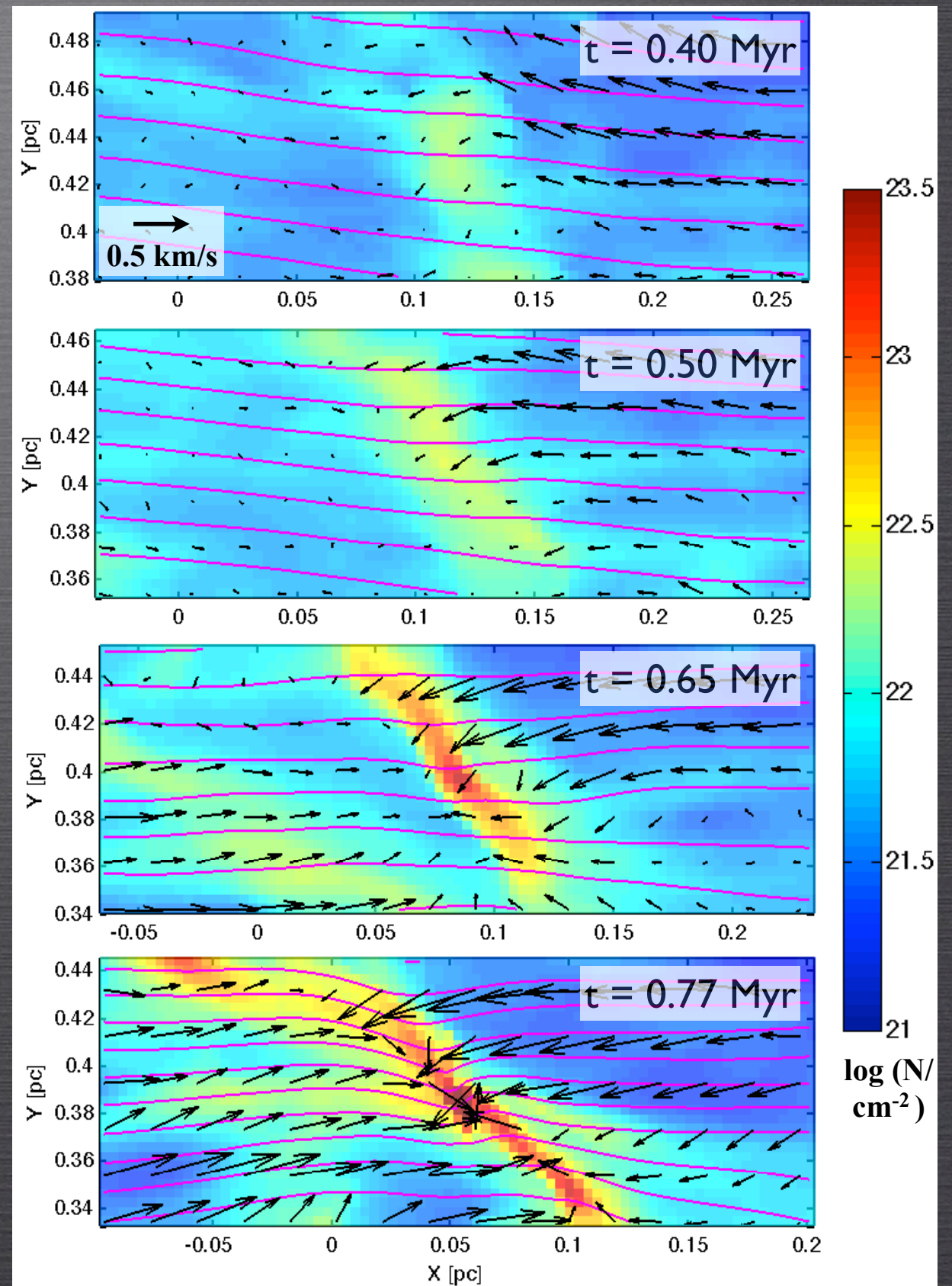
- Models with same pre-shock pressure
⇒ similar core masses, with or without AD



ANISOTROPIC GAS FLOWS

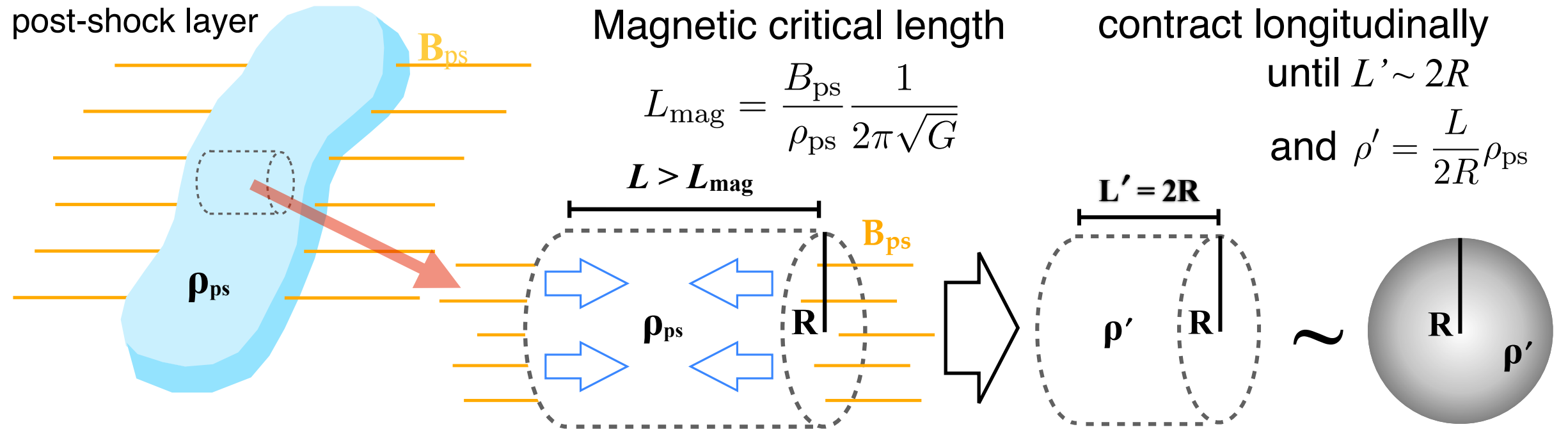


(Chen & Ostriker 2014)



Anisotropic Core Formation

filament formed in post-shock layer



Thermally supercritical: $R \sim R_{\text{BE}}(\rho') = 0.65 \frac{c_s}{\sqrt{G\rho'}} \Rightarrow R = 0.84 \frac{c_s^2}{G\rho_{\text{ps}}L},$

$$M = \pi R^2 L \rho_{\text{ps}} = 2.2 \frac{c_s^4}{G^2 \rho_{\text{ps}} L}$$

Magnetically critical, Anisotropic

$$L = L_{\text{mag}} \Rightarrow \rho_{\text{ps}} L = \frac{B_{\text{ps}}}{2\pi\sqrt{G}}$$

$$M_{\text{crit,cyl}} = 14 \frac{c_s^4}{\sqrt{G^3} B_{\text{ps}}}$$

$$= 1.3 \text{ M}_{\odot} \left(\frac{B_{\text{ps}}}{50 \mu\text{G}} \right)^{-1} \left(\frac{T}{10 \text{ K}} \right)^2$$

$$= 2.8 \frac{c_s}{\sqrt{G^3 \rho_0 v_0^2}} \propto \mathcal{M}^{-1}$$

Magnetically critical, Isotropic

Spherical core with $R = \frac{L_{\text{mag}}}{2} = \frac{B_{\text{ps}}}{4\pi\sqrt{G}\rho_{\text{ps}}}$

$$M_{\text{crit,sph}} = \frac{4\pi R^3}{3} \rho_{\text{ps}} = \frac{1}{48\pi^2 G^{3/2}} \frac{B_{\text{ps}}^3}{\rho_{\text{ps}}^2}$$

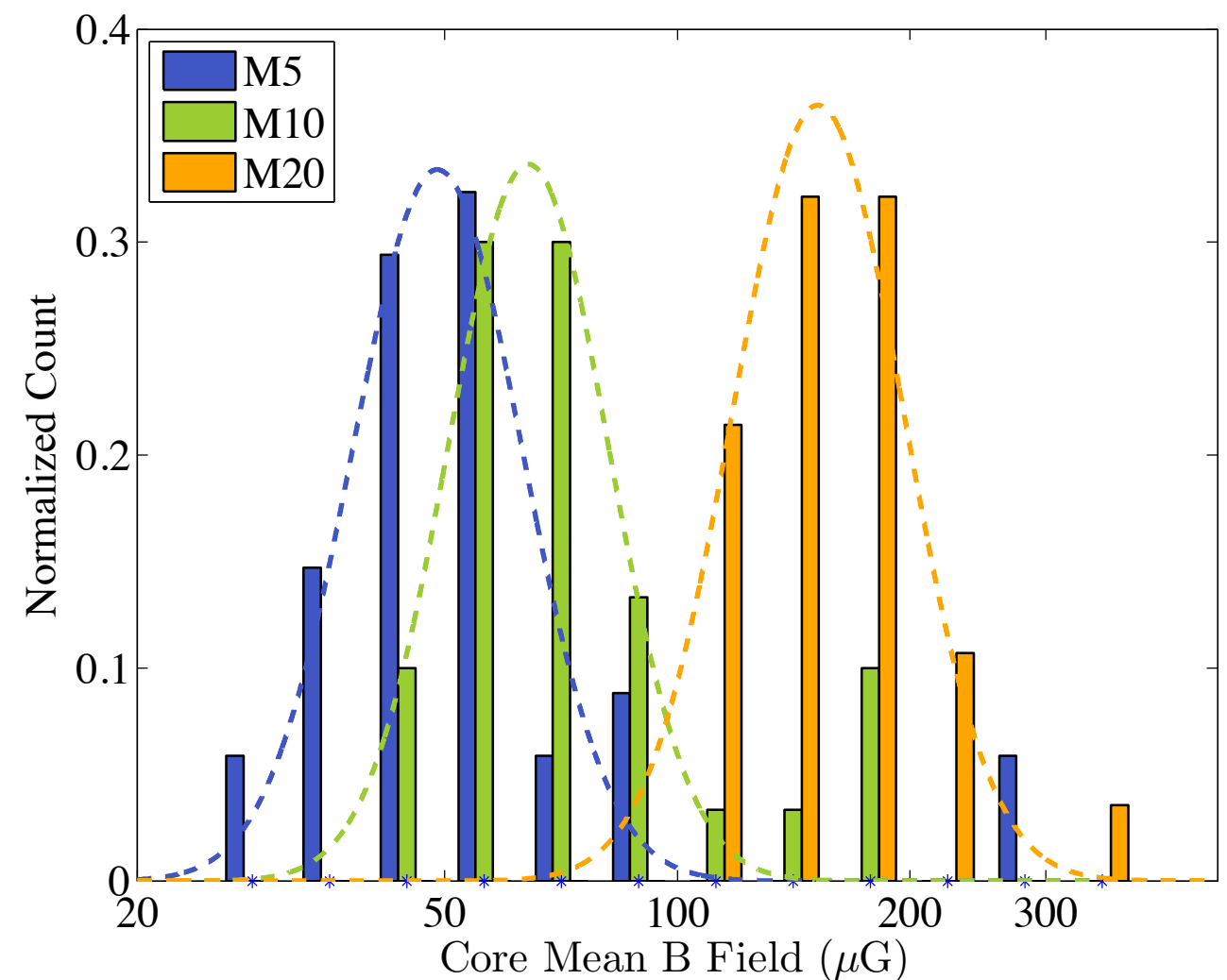
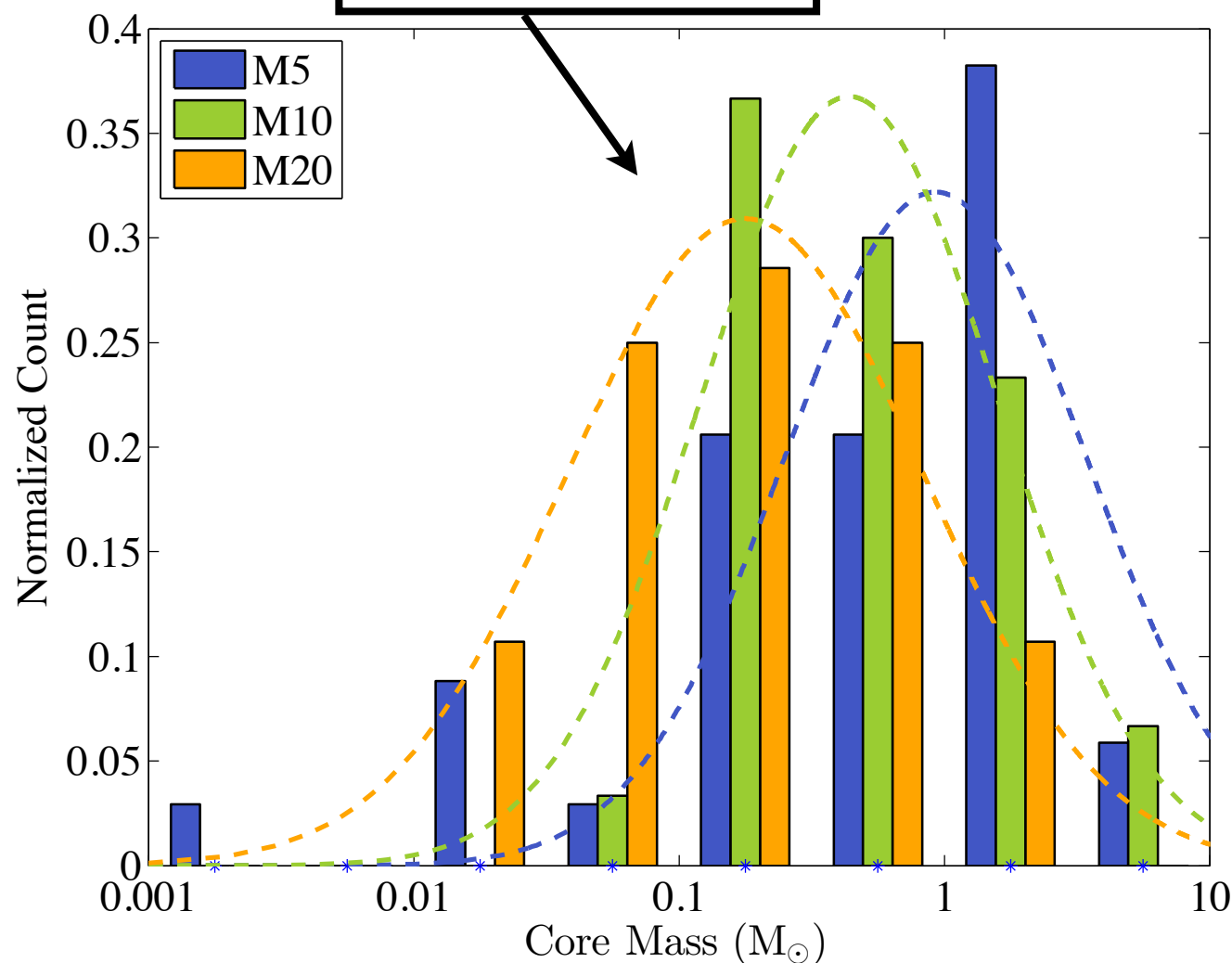
$$= 5.21 \text{ M}_{\odot} \left(\frac{B_{\text{ps}}}{50 \mu\text{G}} \right)^3 \left(\frac{n_{\text{ps}}}{10^4 \text{ cm}^{-3}} \right)^{-2}$$

CORE PROPERTIES

- Core mass varies with inflow Mach number

$$M_{\text{core}} \propto \mathcal{M}^{-0.9}$$

(Chen & Ostriker *in prep.*)



- B_{core} is within a factor of 2 of $B_{\text{post-shock}}$

SUMMARY

- Filament network and core properties found in our simulations are comparable to observations
- Filament transverse velocity gradients provide evidence of condensation in flattened structures
- Magnetically-supercritical, low-mass cores form anisotropically via contraction along $\vec{B}$
- These cores have masses and magnetic fields that depend on pre-shock ρv^2 in cloud